

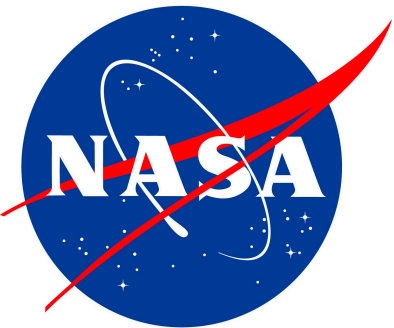
Advanced Modeling & Simulation (AMS) Seminar Series

NASA Ames Research Center, February 2nd, 2021

High-order discontinuous Galerkin method for entry simulations into dusty atmospheric environments

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Objective

Summary of research results from NASA-supported ECF award NNX15AU58G on “Advanced Physical Models and Numerical Algorithms to Enable High-Fidelity Aerothermodynamic Simulations of Planetary Entry Vehicles on Emerging Distributed Heterogeneous Computing Architectures”

Overall research goals

- Develop innovative physical models and advanced numerical methods for reliably predicting aerothermodynamic flows that are relevant to hypersonic and atmospheric entry vehicles in dusty flow environments

Objective

- Develop high-order DG solver for high-speed disperse multiphase flows
 - › Shock capturing method, with a focus on predicting surface heating in hypersonic flows
 - › Lagrangian particle method suited for arbitrary curved, high-aspect-ratio elements
- Assess DG method for hypersonic, particle-laden flow applications
 - › Quantitative comparisons with state-of-the-art finite volume solvers (FUN3D, LAURA)
- Apply developed DG framework to hypersonic dusty flows simulations
 - › Parametric sensitivity study
 - › Hypersonic dusty flow over the ExoMars Schiaparelli capsule

Acknowledgments

NASA

- Michael Barnhardt, Grant Palmer, Khalil Bensassi, Peter Gnoffo, Ethiraj Venkatapathy, Amal Sahai, and entire EDL-team

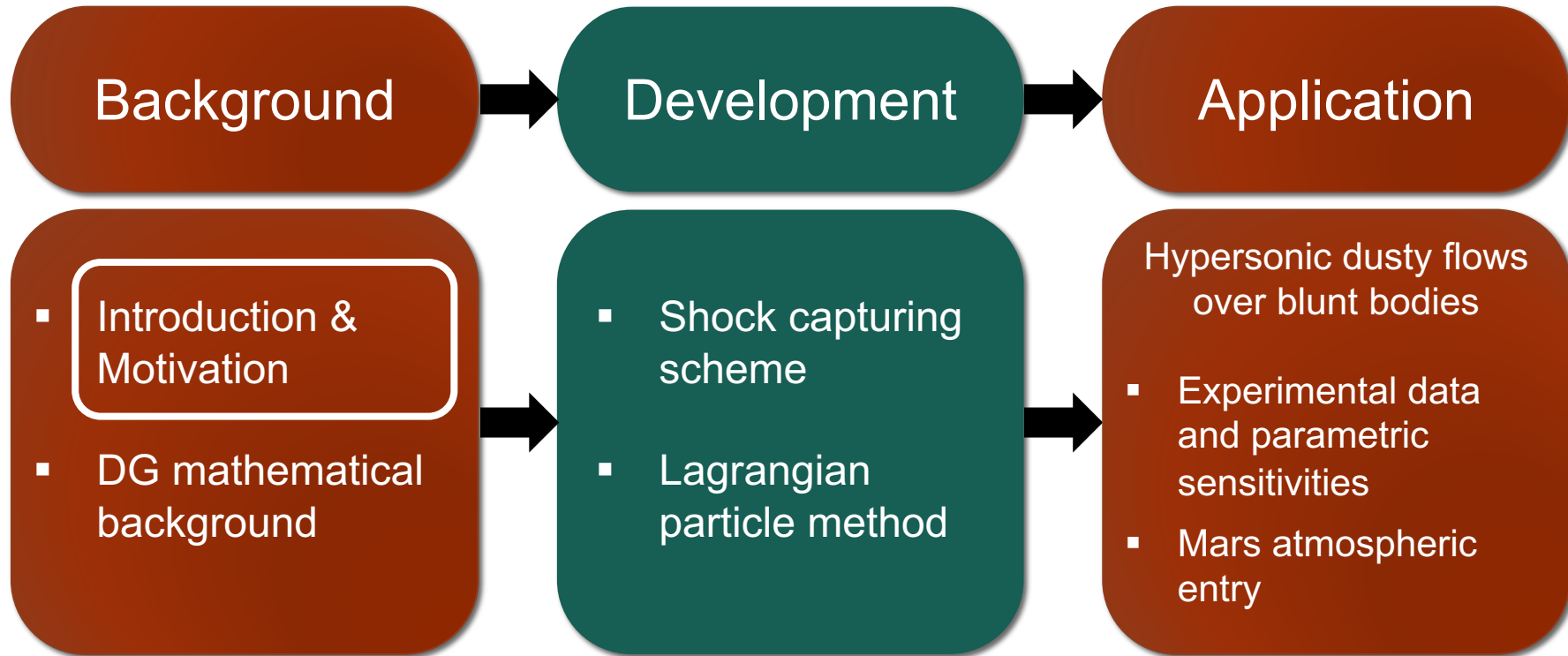
DLR

- Ali Gülhan, Dirk Allofs

Stanford

- Yu Lv, Steven Brill, Narendra Singh, Alex Aiken

Outline

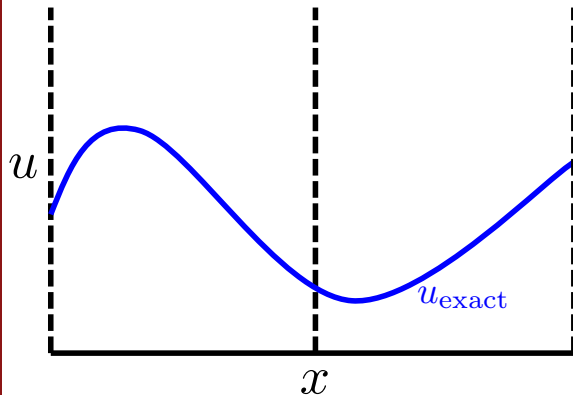


What is DG?

Discontinuous Galerkin methods are a family of numerical methods combining features of finite-element and finite-volume schemes

- Piecewise polynomials used to approximate the solution
 - Arbitrary polynomial order p
 - In each element, solve for polynomial coefficients \tilde{u}

$$u_{\text{approx}}(t, x) = \sum_{i=1}^{p+1} \tilde{u}_i(t) \phi_i(x)$$



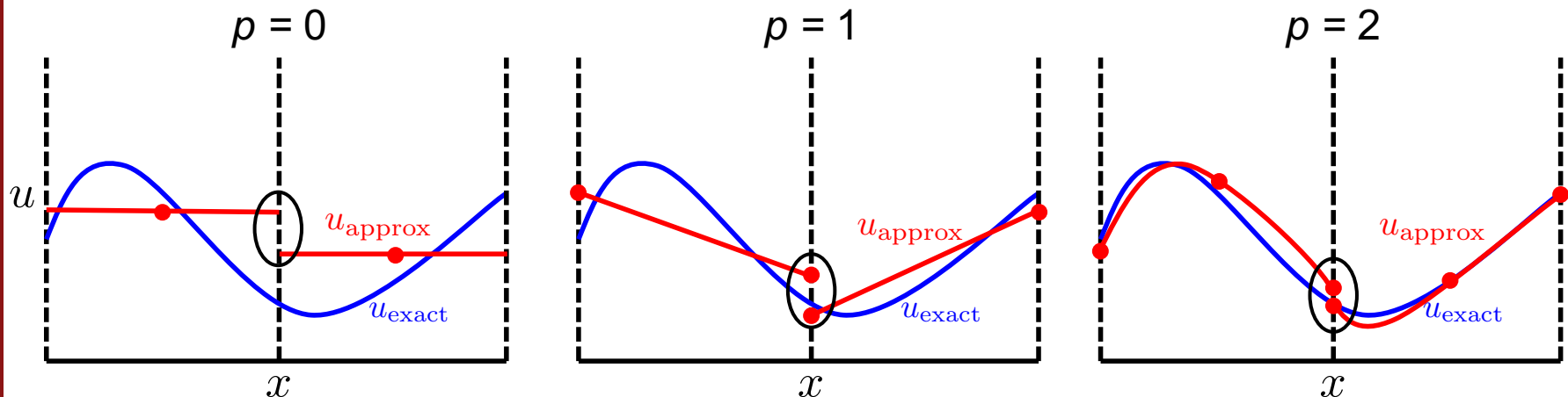
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- Numerical fluxes needed to deal with discontinuities between elements



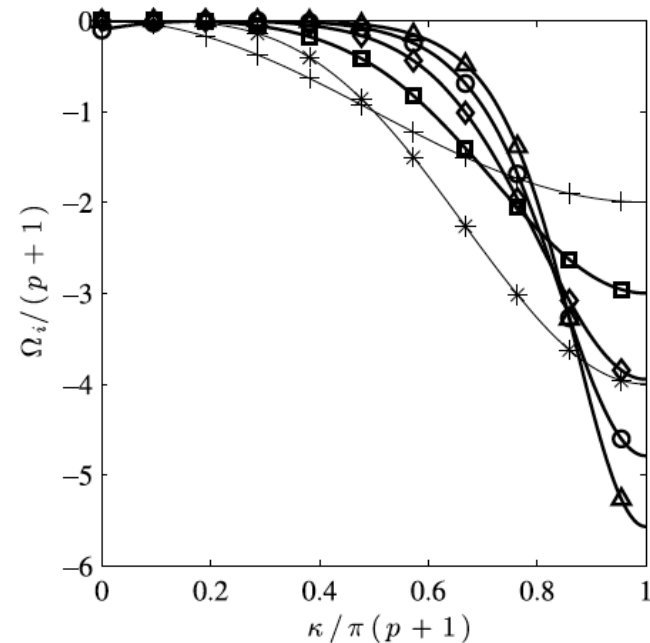
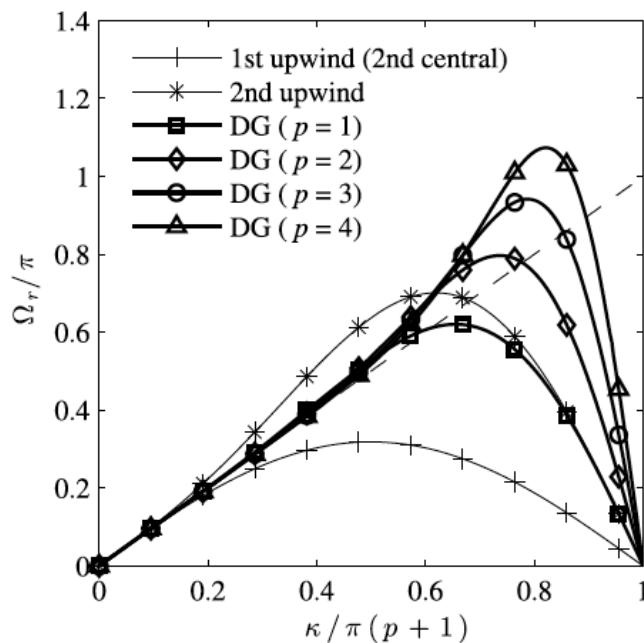
Why DG?

Current state of simulations of multi-physics (turbulent, multiphase, chemically reacting, etc.) flows on complex geometries

- Finite-difference, finite-volume methods
- Low-order (1st or 2nd order)

Potential for high-order DG methods

- Desirable dissipation and dispersion properties



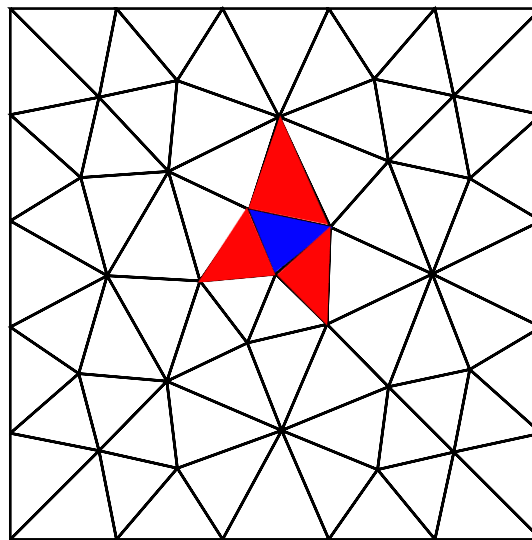
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Potential for high-order DG methods

- Desirable dissipation and dispersion properties
- Suited for HPC (compactness, FLOPs : memory bandwidth)



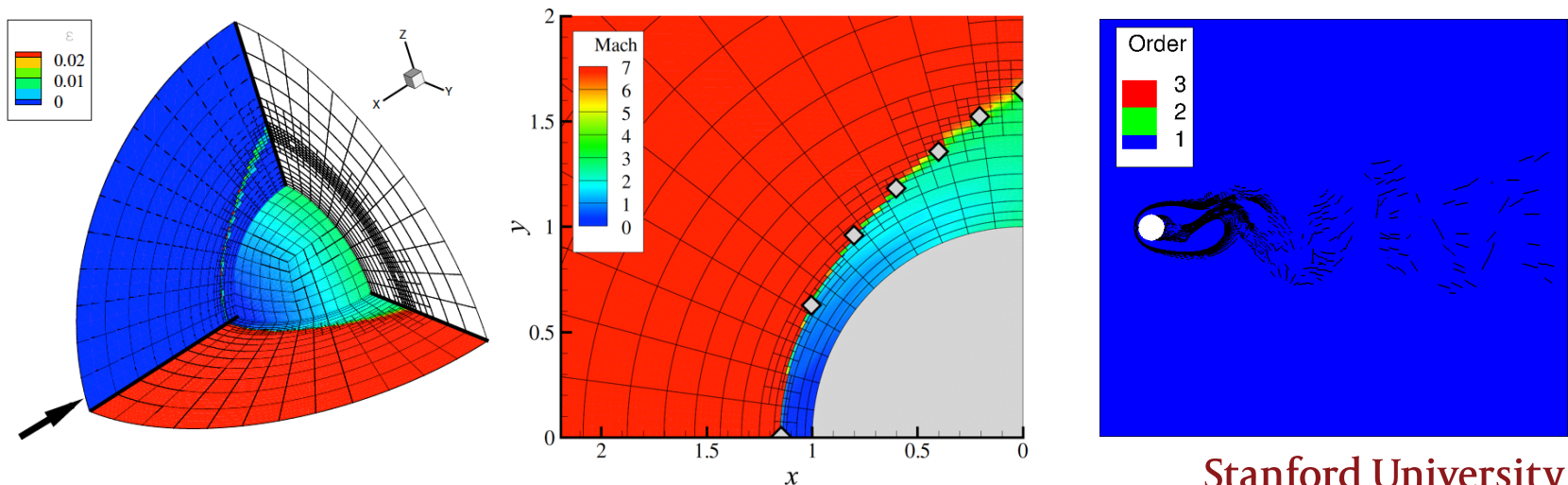
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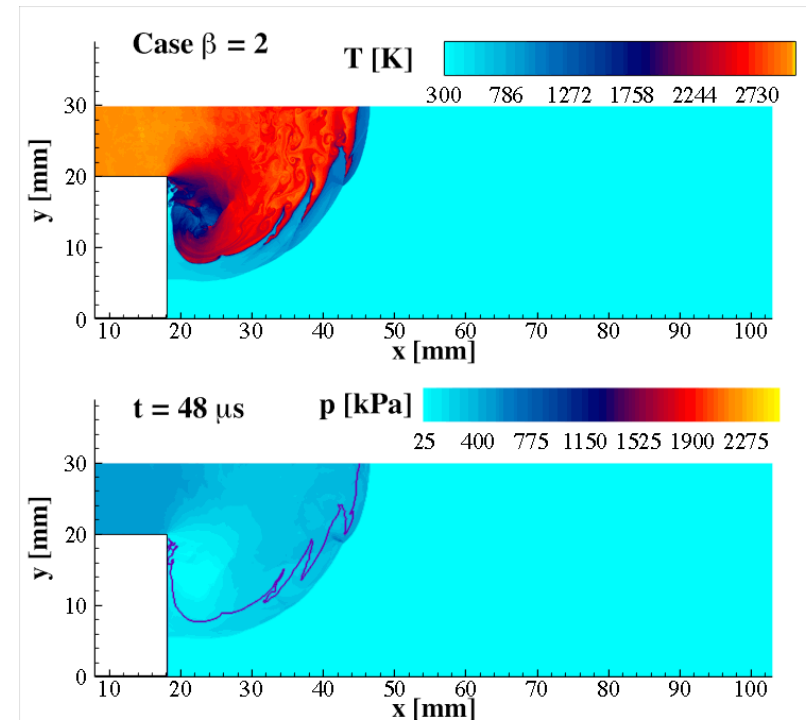
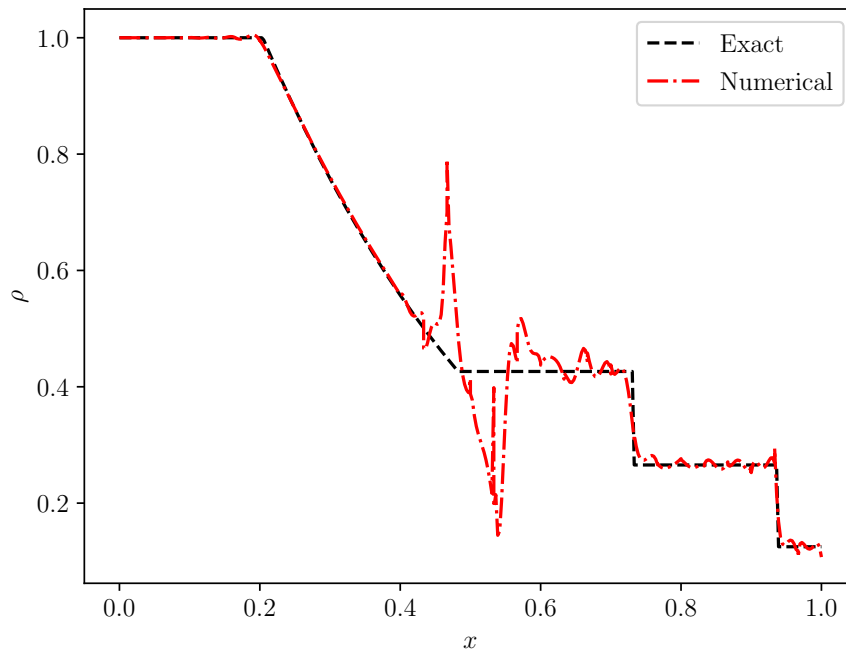
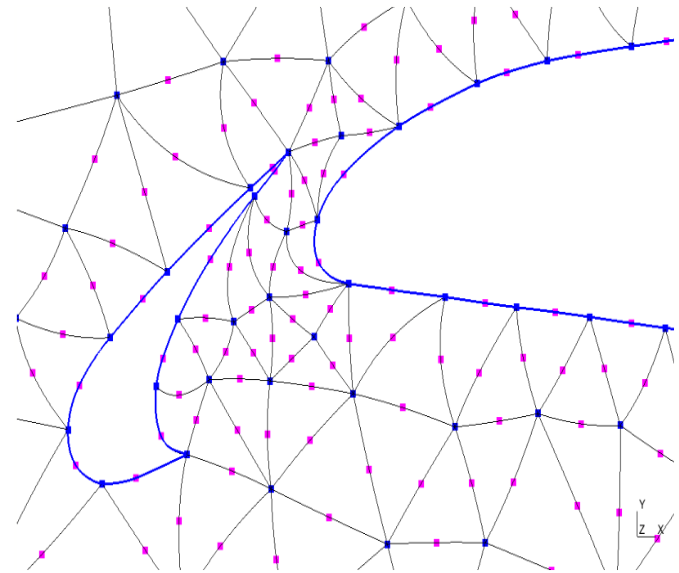
Potential for high-order DG methods

- Desirable dissipation and dispersion properties
- Suited for HPC (compactness, FLOPs : memory bandwidth)
- Local mesh-adaptation (h) and refinement in polynomial order (p)

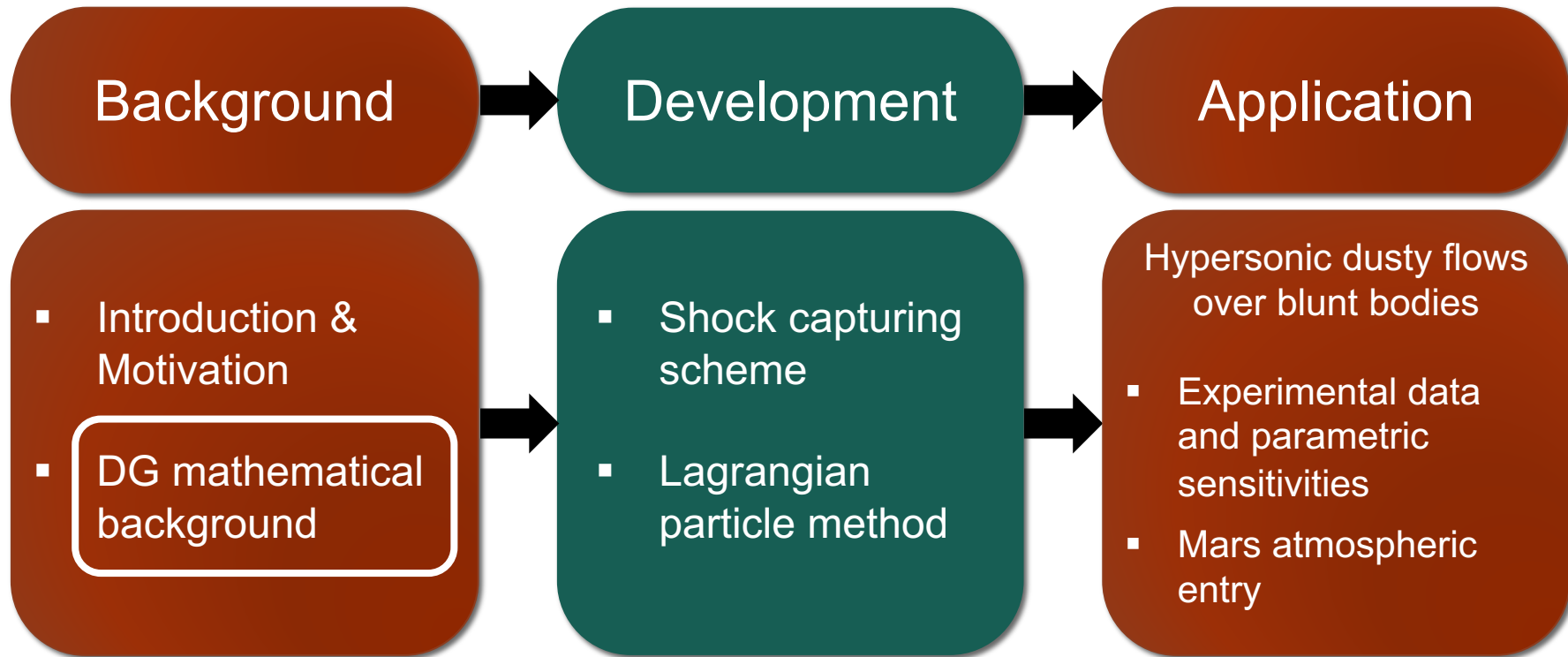


Challenges

- Efficient time stepping
- Efficient automatic *hp*-adaptation
- High-order curved meshes
- Numerical instabilities
- Extensions to complex physics



Outline



Discontinuous Galerkin discretization

Governing equations

$$\partial_t \mathbf{U} + \nabla \cdot \mathbf{F}_s = \nabla \cdot \mathbf{F}_v + \mathbf{S}$$

$$\mathbf{U} = \begin{bmatrix} \rho \\ \rho \mathbf{u} \\ \rho E \end{bmatrix}, \quad \mathbf{F}_s = \begin{bmatrix} \rho \mathbf{u} \\ \rho \mathbf{u} \otimes \mathbf{u} + P \mathbb{I} \\ \mathbf{u}(\rho E + P) \end{bmatrix}, \quad \mathbf{F}_v = \begin{bmatrix} 0 \\ \boldsymbol{\tau} \\ \mathbf{u} \cdot \boldsymbol{\tau} - q \end{bmatrix}$$

Partition computational domain

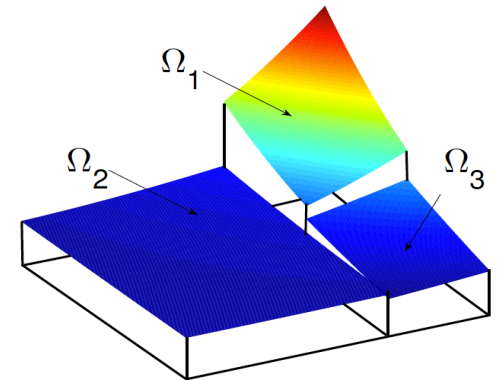
$$\Omega = \bigcup_{e=1}^{N_e} \Omega_e$$

Approximate solution with polynomials of order p

$$\mathbf{U} \approx \mathbf{U}_h = \bigoplus_{e=1}^{N_e} \mathbf{U}_h^e$$

$$\mathbf{U}_h^e(\mathbf{x}, t) = \sum_{n=1}^{N_b} \tilde{\mathbf{U}}_n^e(t) \phi_n(\mathbf{x})$$

Multiply governing equations by ϕ_m and integrate



Discontinuous Galerkin discretization

Find the basis coefficients $\tilde{U}^e(t)$ that discretely satisfy

$$\sum_{n=1}^{N_b} d_t \tilde{U}_n^e(t) \int_{\Omega_e} \phi_m \phi_n d\Omega_e + \underbrace{\int_{\Omega_e} \phi_m \nabla \cdot \mathbf{F}_s d\Omega_e}_{\text{HLLC}^1, \text{Roe}^2, \text{AUSM}^+, \text{SLAU}^4} = \underbrace{\int_{\Omega_e} \phi_m \nabla \cdot \mathbf{F}_v d\Omega_e + \int_{\Omega_e} \phi_m \mathbf{S} d\Omega_e}_{\text{Interior Penalty}^5, \text{BR2}^6} \quad \forall \phi_m \quad m = 1, \dots, N_b$$

HLLC¹

Roe²

AUSM⁺³

SLAU⁴

Interior Penalty⁵

BR2⁶

Evaluate integrals using Gaussian quadrature

[1] Toro, et al., SW, 1994;

[2] Roe, JCP, 1981;

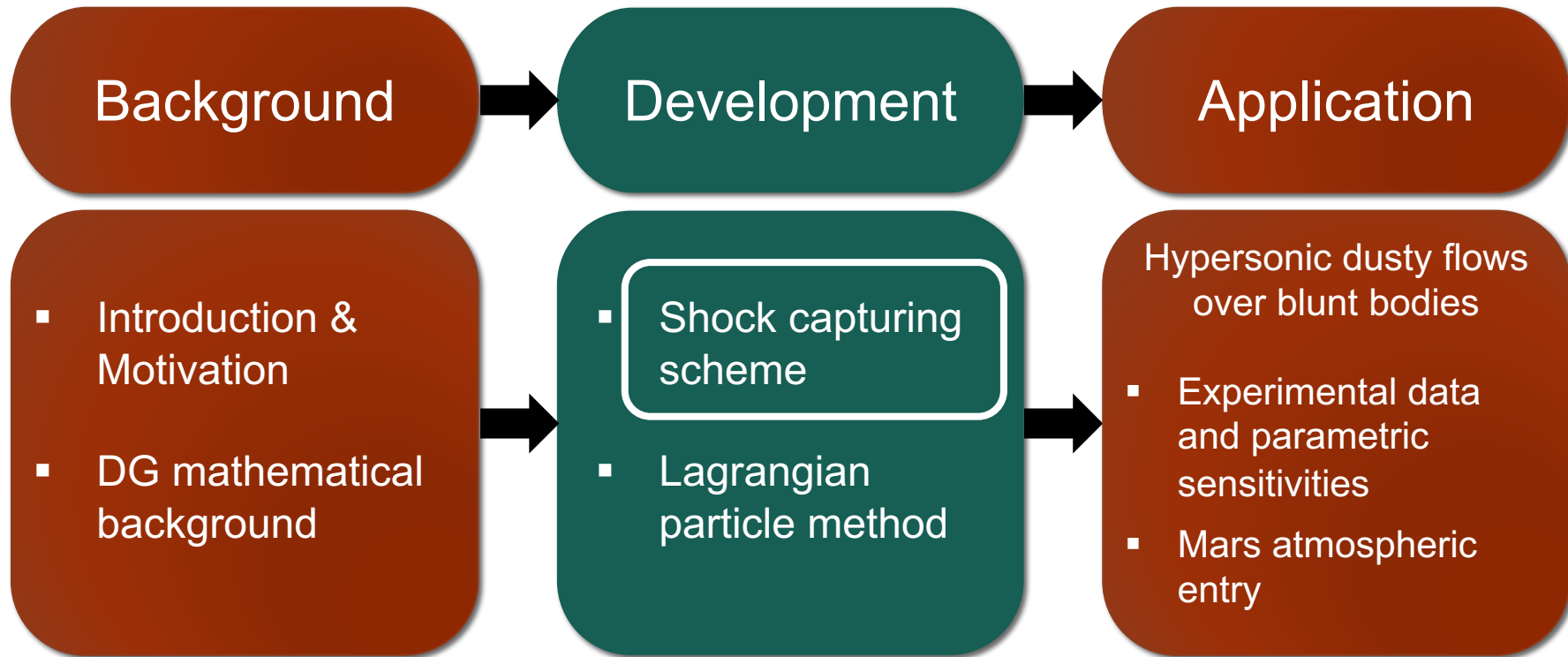
[3] Liou, JCP, 1996;

[4] Shima and Kitamura, AIAA J., 2011;

[5] Arnold, et al. SIAM JNA, 1981;

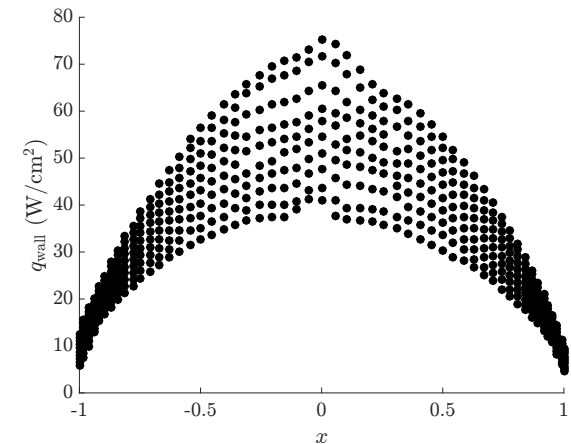
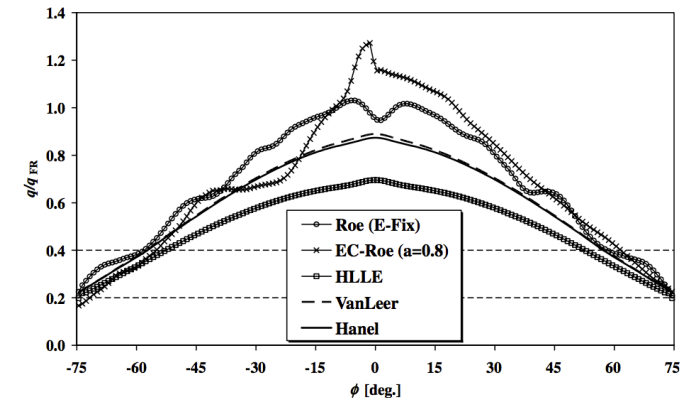
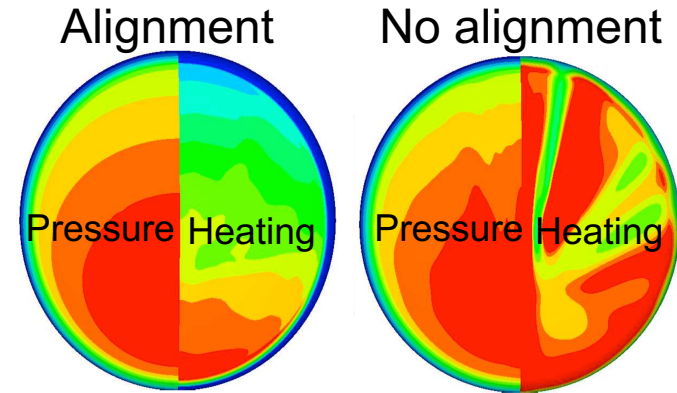
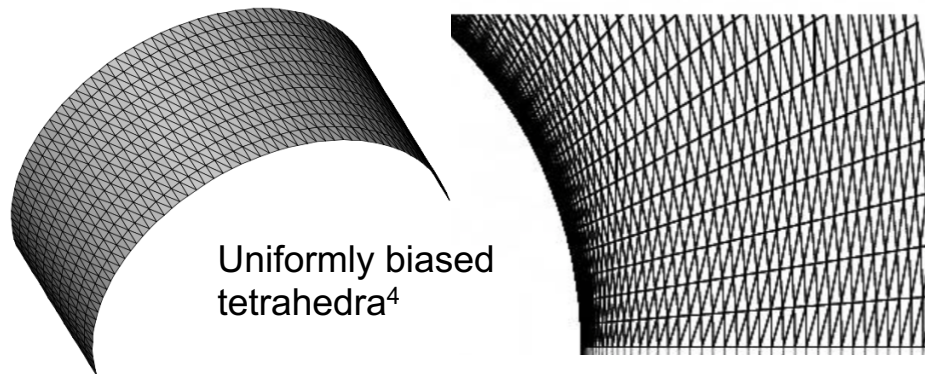
[6] Bassi & Rebay, JCP, 1997.

Outline



Background: finite-volume

- Surface heating predictions extremely sensitive to:
 - Mesh-shock alignment¹
 - Inviscid flux function^{2,3}
 - Limiter, parameters, etc.^{2,3}
- Gnoffo computed hypersonic flow over a cylinder using **uniformly biased tetrahedra**⁴



[1] Candler et al., AIAA Paper, 2009
[2] Kitamura et al., AIAA Journal, 2010

[3] Kitamura, AIAA Paper, 2013
[4] Gnoffo and White, AIAA Paper, 2004

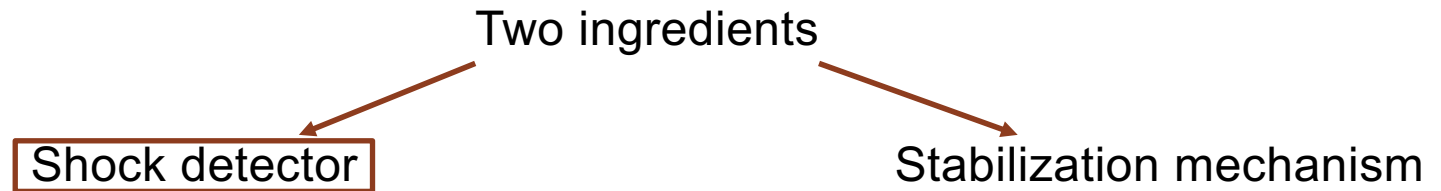
Background: DG

- Robust shock capturing is still an active area of research
- Limited work on using high-order DG to compute viscous hypersonic flows^{1,2,3,4}

Goals:

- Develop a simple and robust shock capturing method
- Assess sensitivities of DG heating predictions to mesh-shock alignment and choice of inviscid and viscous flux functions
- Perform quantitative comparisons with FUN3D and LAURA

Shock capturing method



$$S_e = \sqrt{\frac{\int_{\Omega_e} (s/\bar{s} - 1)^2 d\Omega}{|\Omega_e|}}$$

Intra-element variations

- Simple
- Compact
- Large near discontinuities, small in smooth regions

Identify **troubled elements**

Shock capturing method



Artificial viscosity (AV) term $\nabla \cdot \mathbf{F}_{\text{AV}}$

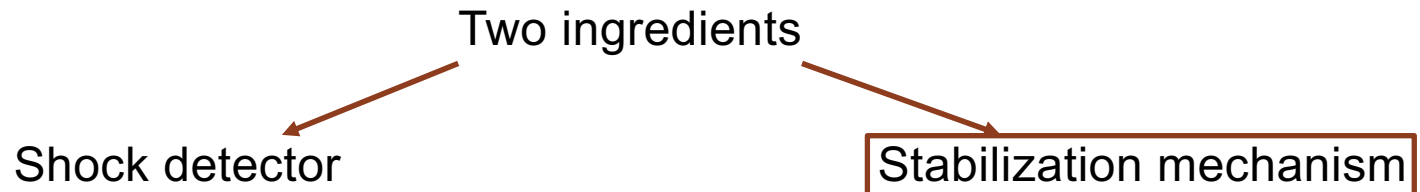
$$\mathbf{F}_{\text{AV}} = \mathbf{H} : \nabla \mathbf{U}$$

$$\mathbf{H}_{ik} = \eta_{\text{AV}} \frac{h_{ik}}{\bar{h}} \mathbb{I}$$

Add initially elementwise-constant AV to troubled elements

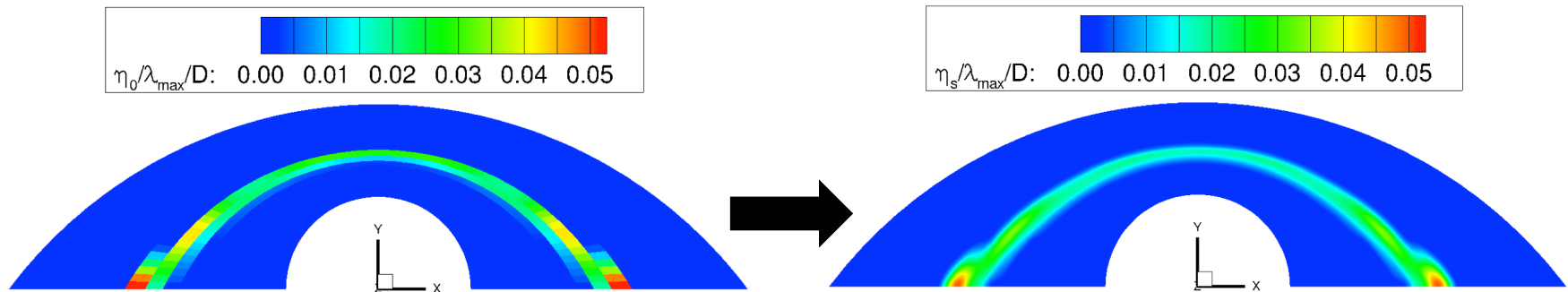
$$\eta_{\text{AV},0}^e = C_\eta \lambda_{\text{max}} \frac{\bar{h}}{\max(1, p)} \bar{S}_e$$

Shock capturing method



Smooth AV

$$\eta_{AV,s} - \nabla \cdot [(C_\Delta \Delta)^2 \nabla \eta_{AV,s}] = \eta_{AV,0}$$



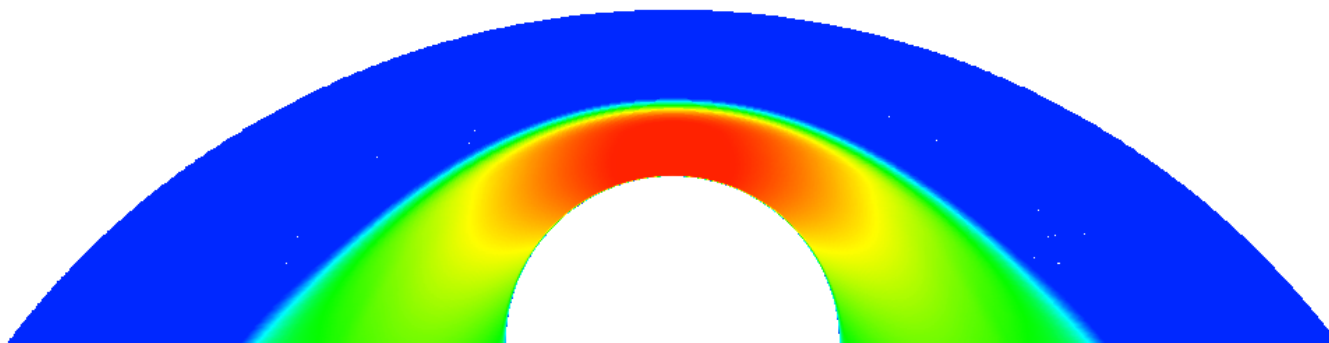
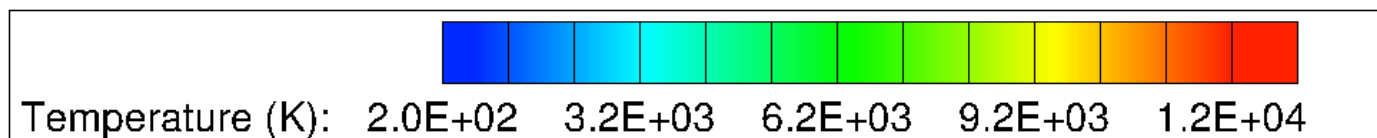
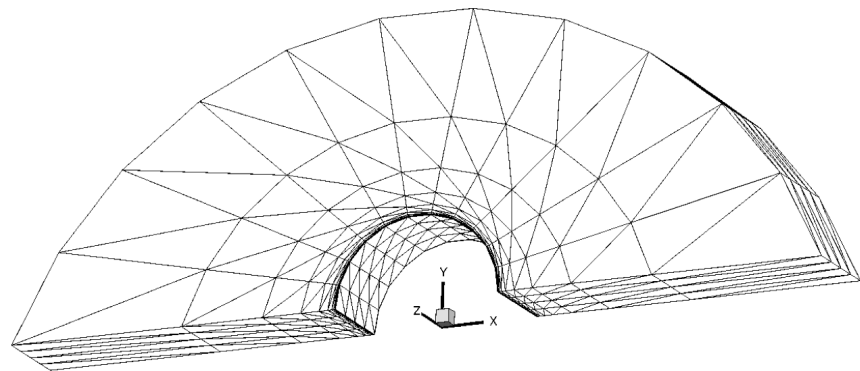
Reduce AV in boundary layers

- Homogeneous Dirichlet BCs at no-slip walls
- Pass $\eta_{AV,s}$ through sinusoidal filter² to get η_{AV}

Test case: Mach 17.6 flow over circular half-cylinder

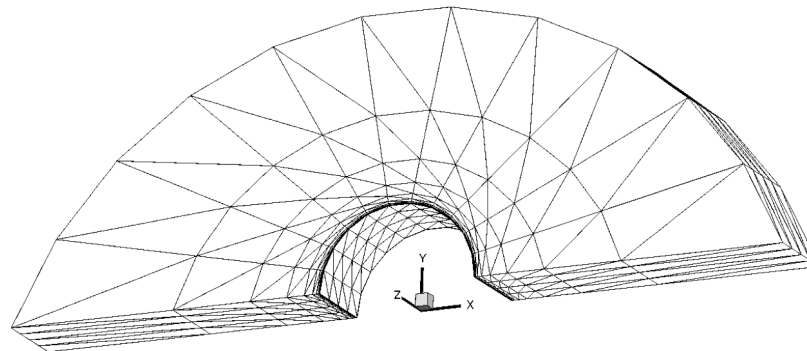
Uniformly biased tetrahedral mesh

- DG: $(p+1)$ -order accuracy, where $p = 1, 2, 3$
- Quantitative comparisons with FUN3D
 - 2nd-order accuracy
- Influences of inviscid and viscous flux functions on heating predictions

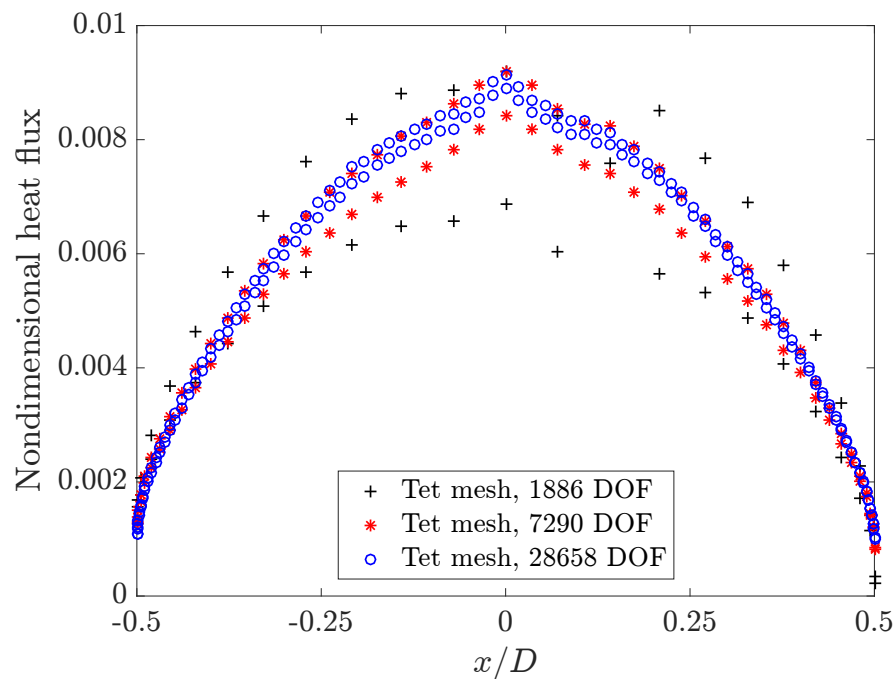


Ma	p -order
17.6	1-3

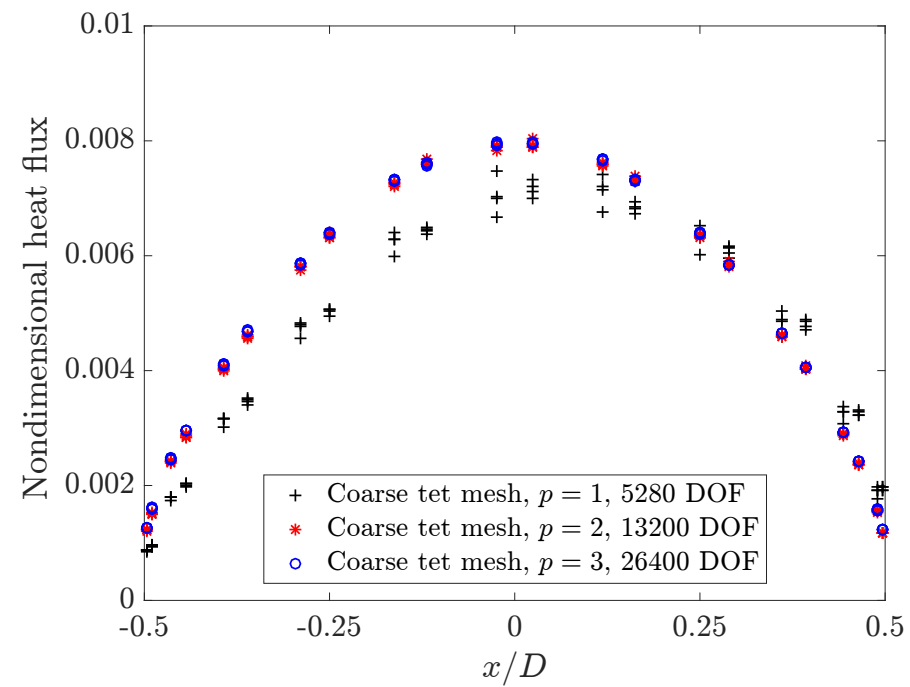
FUN3D heat flux comparisons



FUN3D

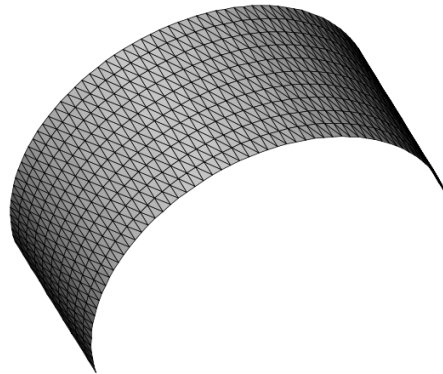


DG

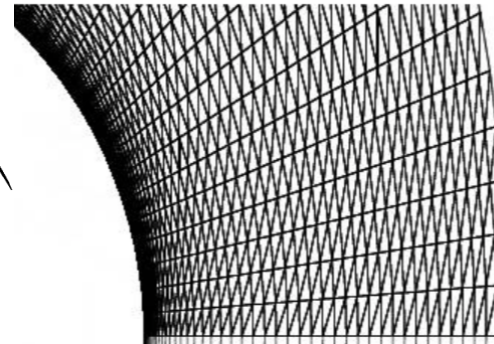
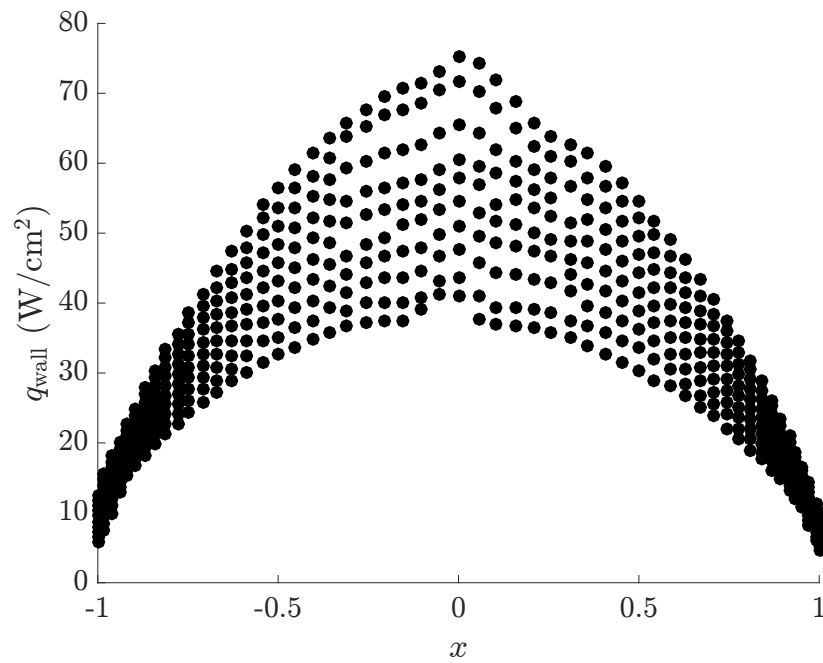


Ma	p -order
17.6	1-3

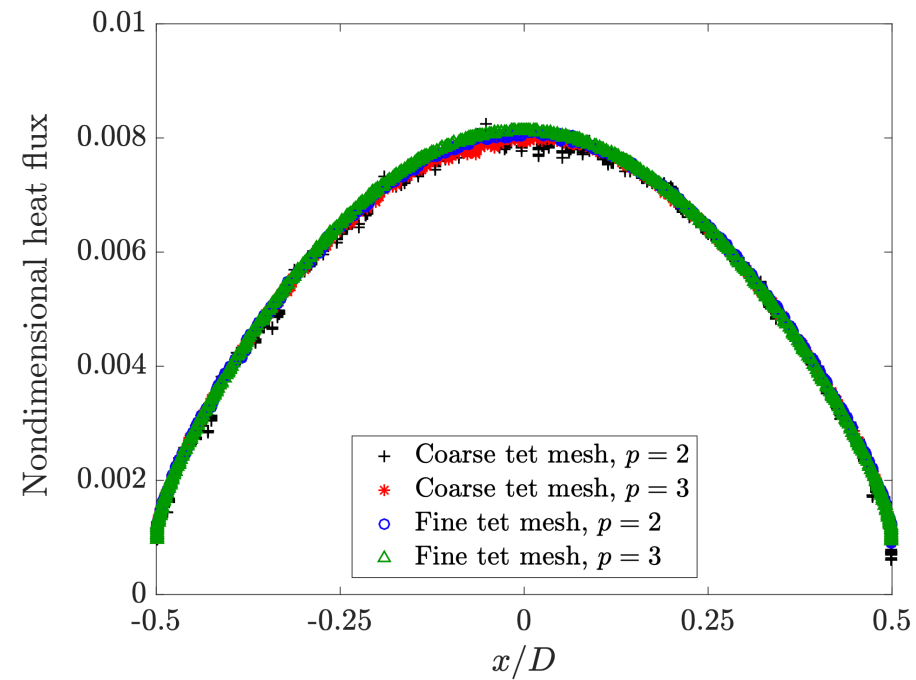
FUN3D heat flux comparisons



FUN3D

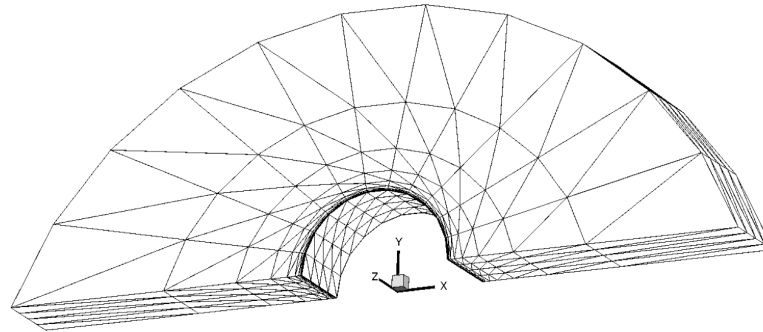


DG

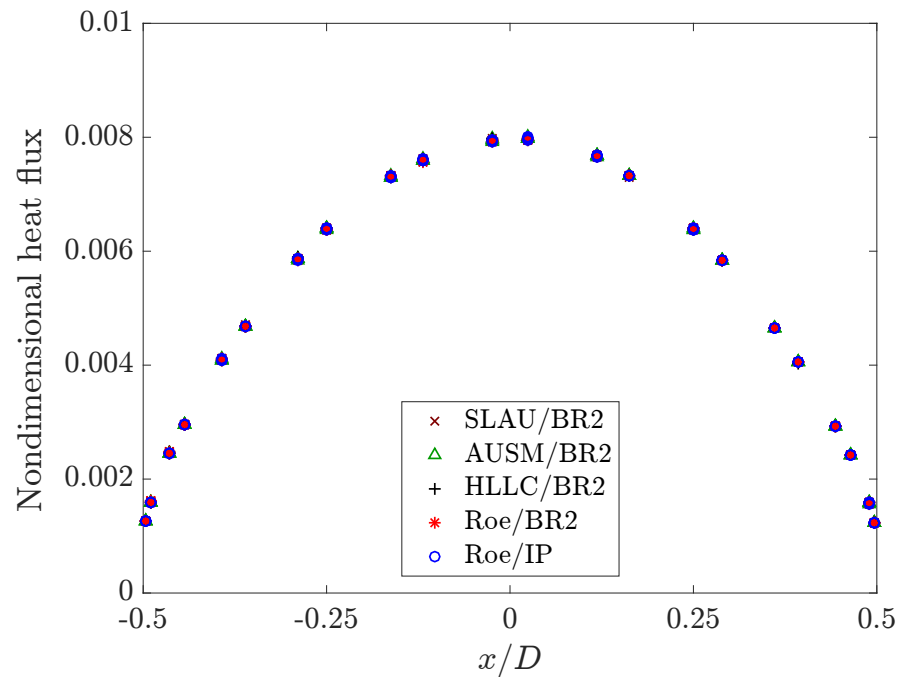


Ma	p -order
17.6	3

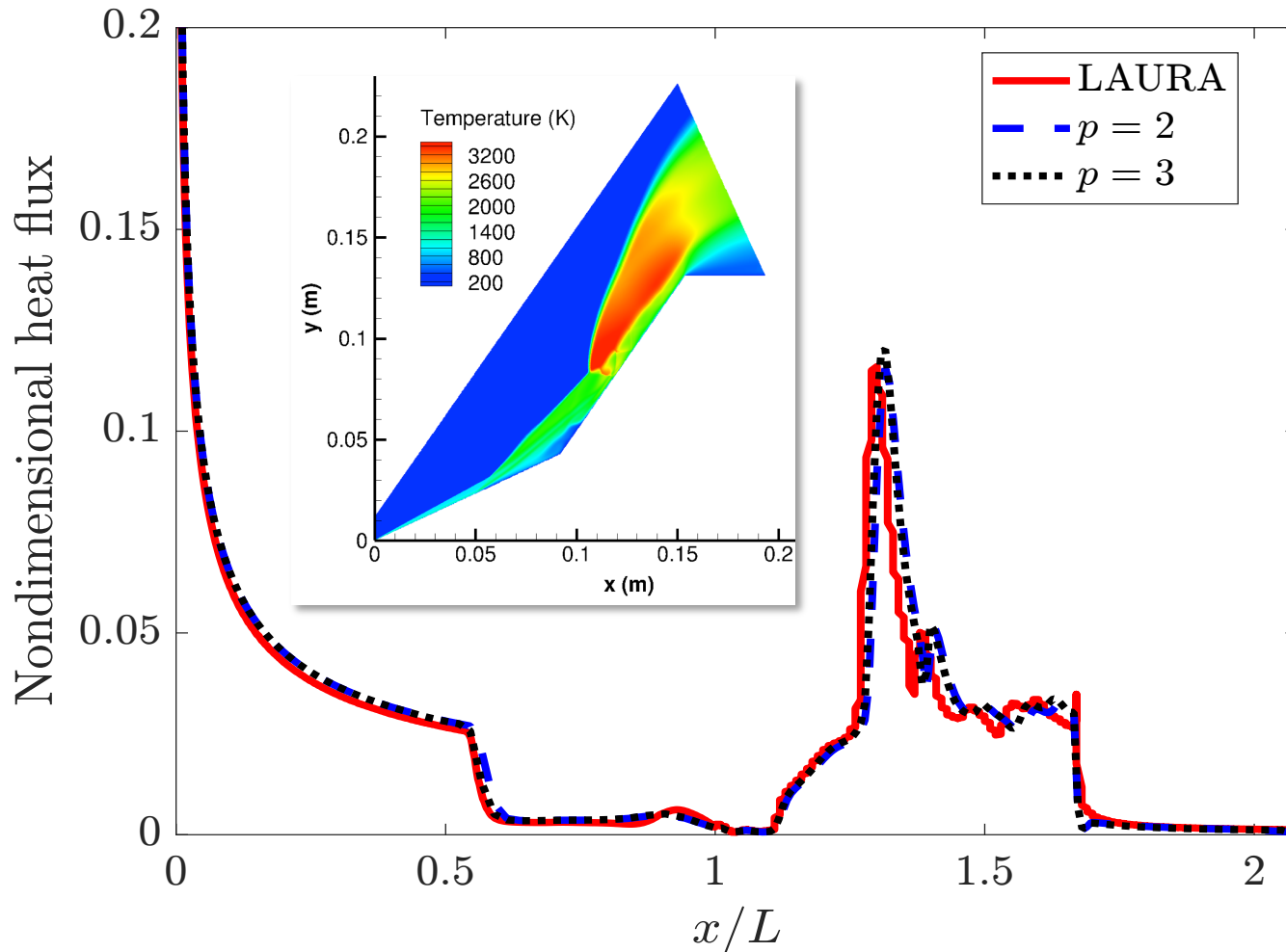
Sensitivity to inviscid and viscous flux functions



Coarse tet mesh, $p = 3$



Test case: hypersonic flow over double cone

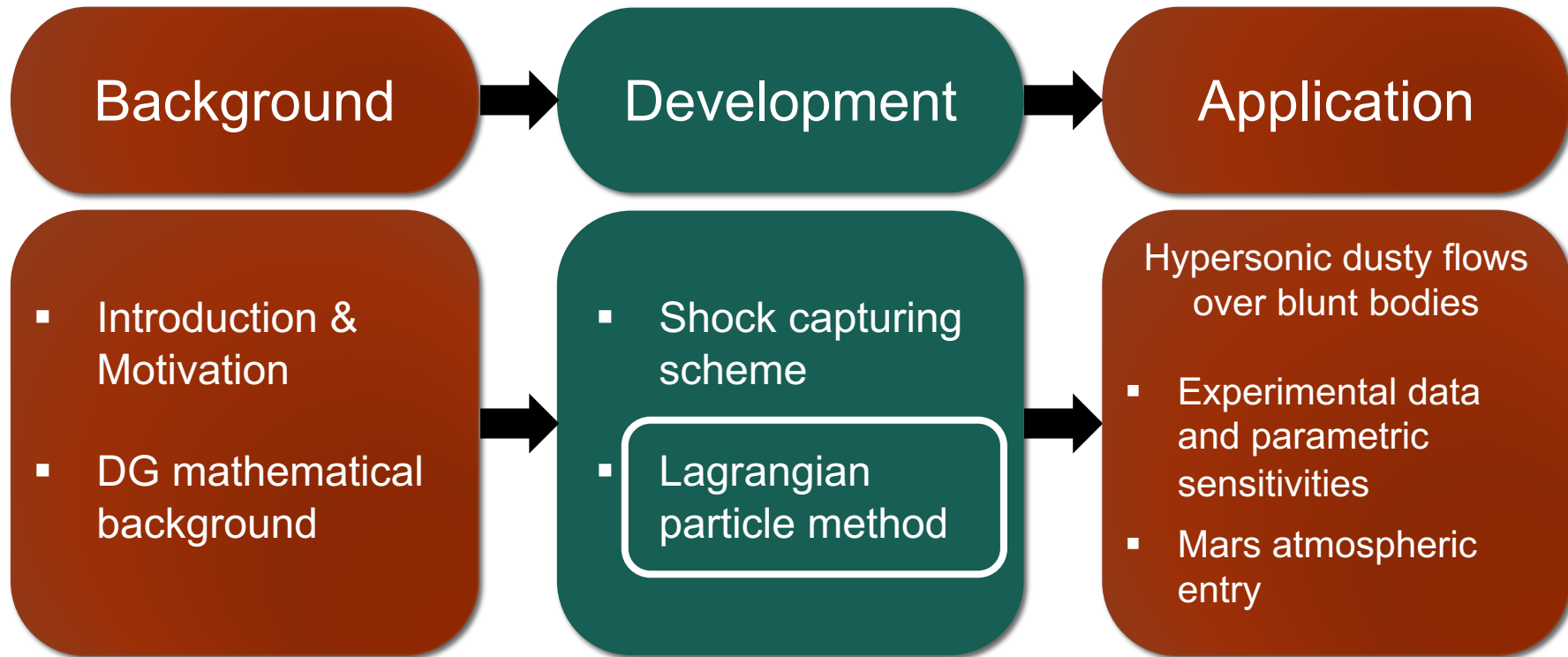


→ DG solution reaches convergence with ~half # DOF of LAURA solution

Summary

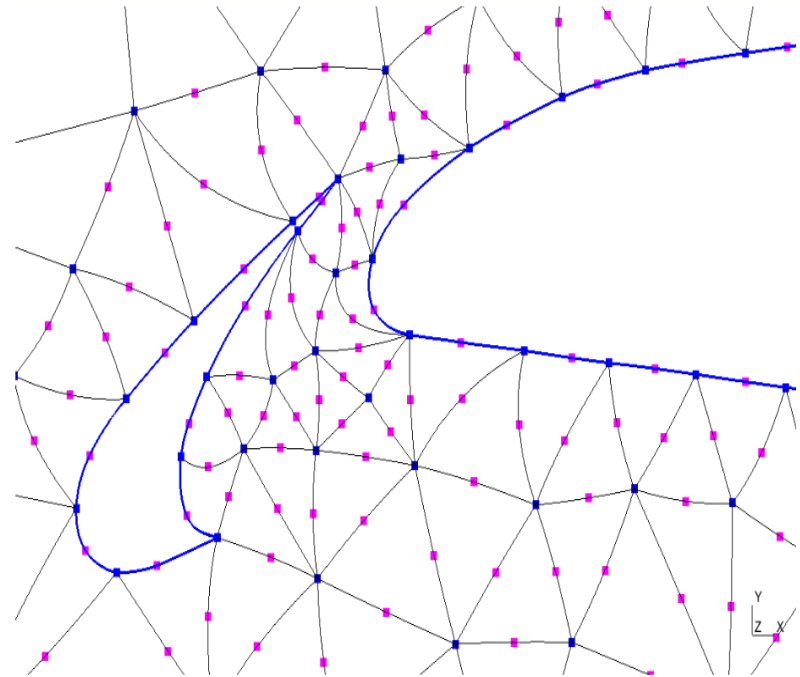
- Developed a simple and robust robust shock capturing method for DG
- Applied the proposed formulation to viscous hypersonic flows
- DG heating predictions are much less sensitive to mesh-shock alignment and choice of inviscid flux function than FV heating predictions
- Fewer degrees of freedom are required to achieve mesh convergence in DG solutions

Outline



Background: particle-laden flows

- Several point-particle Euler-Lagrange formulations that rely on high-order numerical methods
 - Finite-difference^{1,2}
 - Spectral methods^{3,4}
- Limited development in the context of:
 - DG^{5,6,7}
 - Curved elements



[1] Jacobs and Don, *J Comp Phys*, 2008.

[2] Jacobs et al., *Theor Comp Fluid Dyn*, 2012

[3] Bagchi and Balachandar, *J Fluid Mech*, 2002

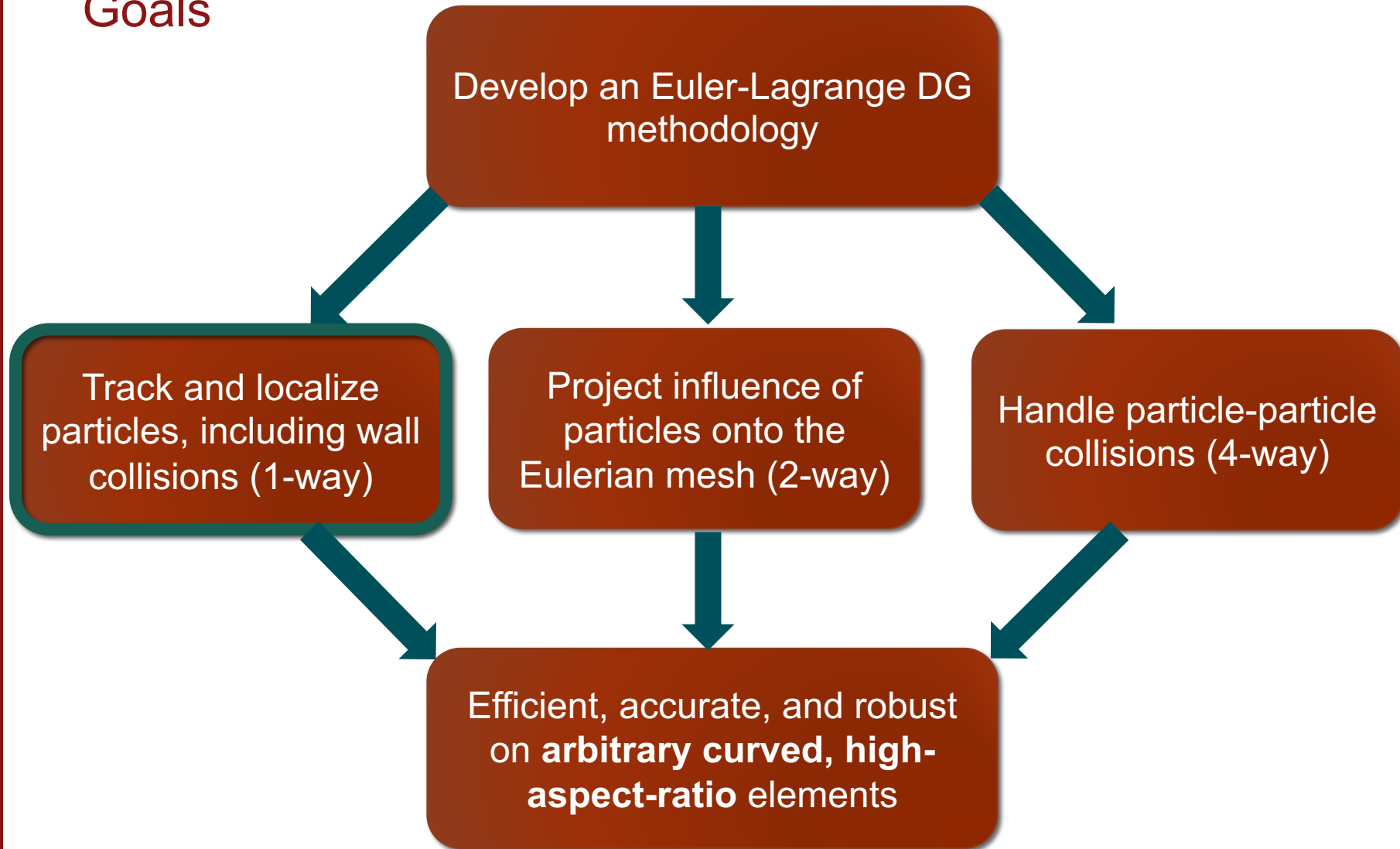
[4] Akiki et al., *J Comp Phys*, 2017

[5] Sengupta et al., *Int J Multiph Flow*, 2009

[6] Huang et al, *Comput Methods Appl Mech Eng*, 2019

[7] Zwick and Balachandar, *Int J High Perform Comput Appl*, 2020

Goals



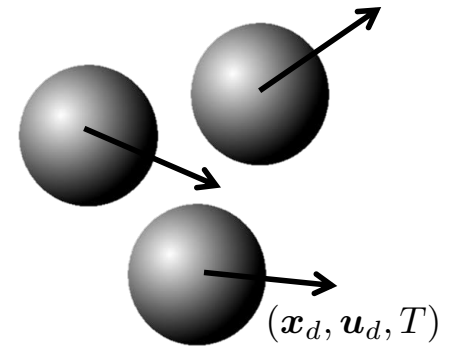
Lagrangian Particle Tracking

- Individual particles (“d”) are tracked in the flow field as point sources
- Carrier gas (“c”) described with Eulerian field

Position: $\frac{d\mathbf{x}_d}{dt} = \mathbf{u}_d$

Momentum: $m_d \frac{d\mathbf{u}_d}{dt} = \mathbf{F} = \mathbf{F}_{qs} + \mathbf{F}_{\text{other}}$

Energy: $m_d c_d \frac{dT_d}{dt} = Q = Q_{qs} + Q_{\text{other}}$



- Quasi-steady contributions

Drag: $\mathbf{F}_{qs} = \frac{1}{8} \pi D^2 \rho_c (\mathbf{u}_c - \mathbf{u}_d) |\mathbf{u}_c - \mathbf{u}_d| C_D$

Heating Rate: $Q_{qs} = \pi D \kappa_c (T_c - T_d) \text{Nu}$

$$C_D = f(\text{Re}_d, \text{Ma}_d)$$

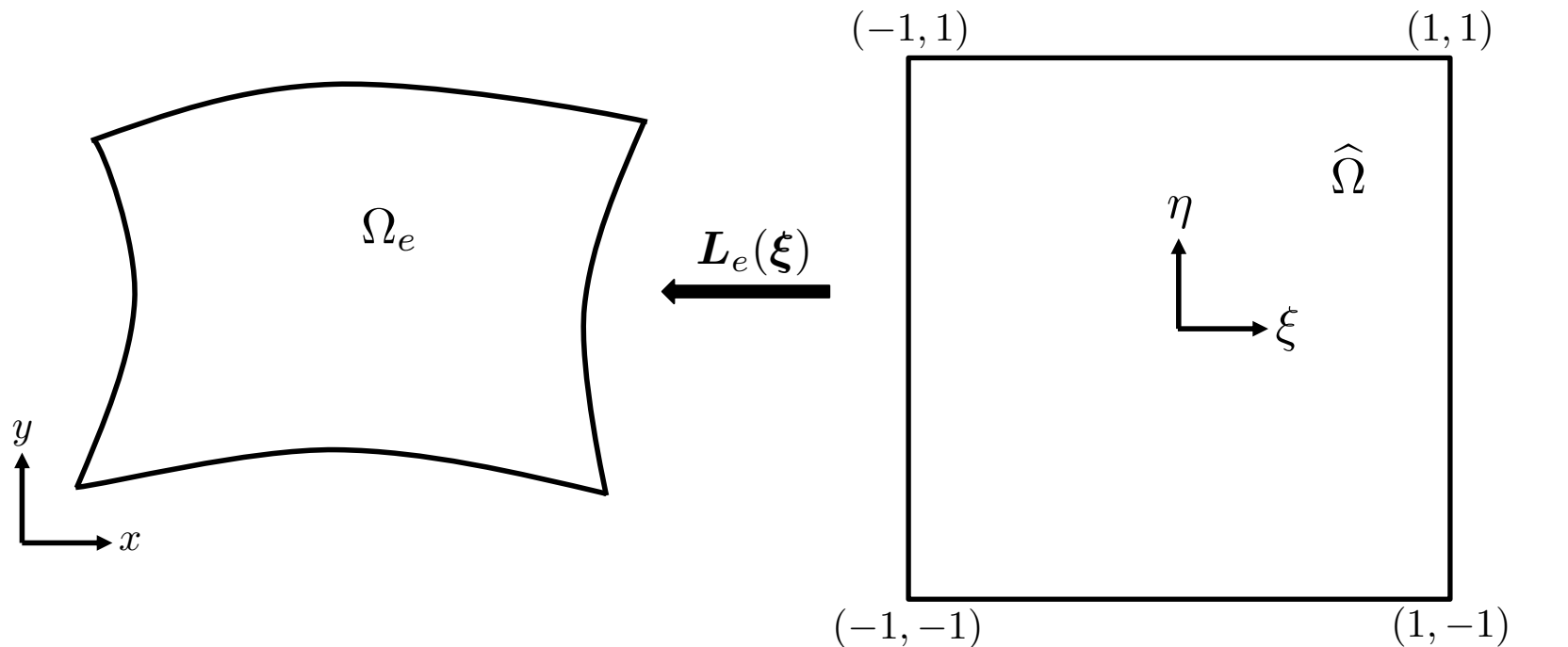
$$\text{Nu} = f(\text{Re}_d, \text{Ma}_d)$$

Particle search-locate procedure

Algorithm by Allievi and Bermejo

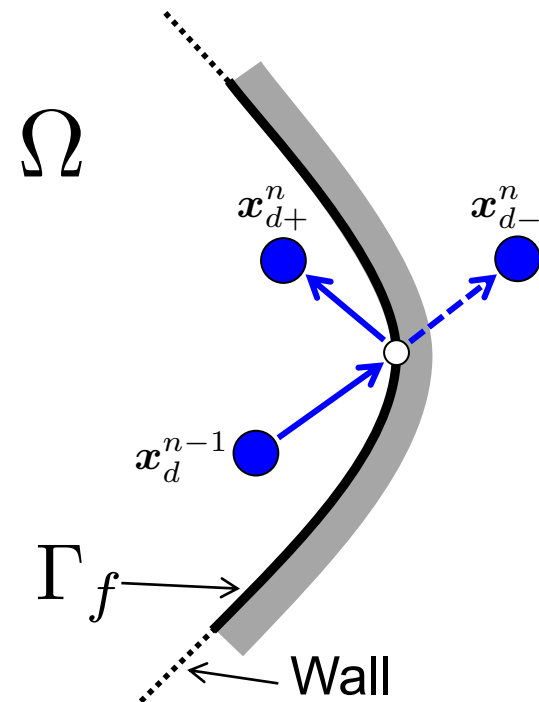
- Compatible with **arbitrary curved elements**
- Relies on the geometric mapping

$$\mathbf{L}_e(\boldsymbol{\xi}) \equiv \begin{bmatrix} x \\ y \end{bmatrix} = \sum_{i=1}^{N_n} \psi_i(\boldsymbol{\xi}) \begin{bmatrix} x_i^e \\ y_i^e \end{bmatrix}$$



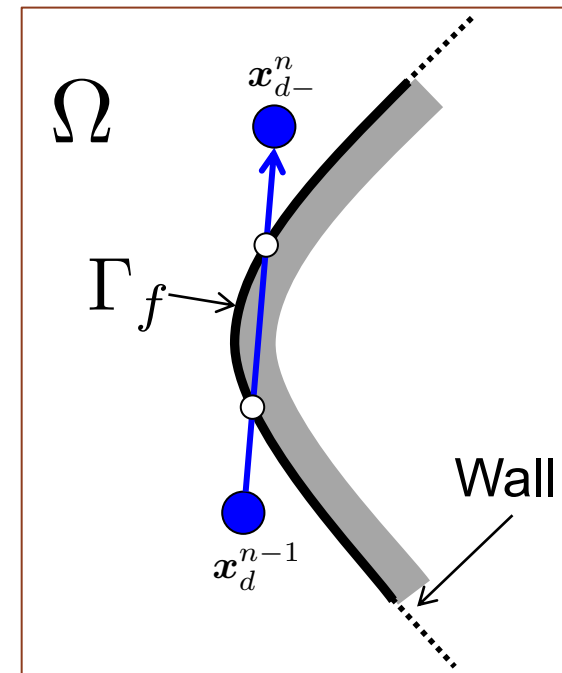
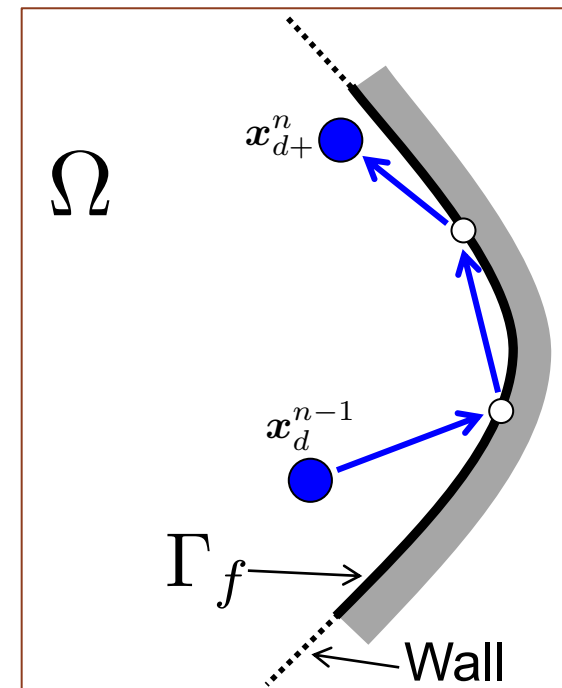
Particle-wall collisions

- Extended the search-locate procedure to account for hard-sphere particle-wall collisions
- Apply Newton search to compute intersection between particle trajectory and boundary
- **Curved, high-aspect-ratio** elements
 - Pathological cases¹



Particle-wall collisions

- Extended the search-locate procedure to account for hard-sphere particle-wall collisions
- Apply Newton search to compute intersection between particle trajectory and boundary
- **Curved, high-aspect-ratio elements**
 - Pathological cases¹
 - Finite size of particles²

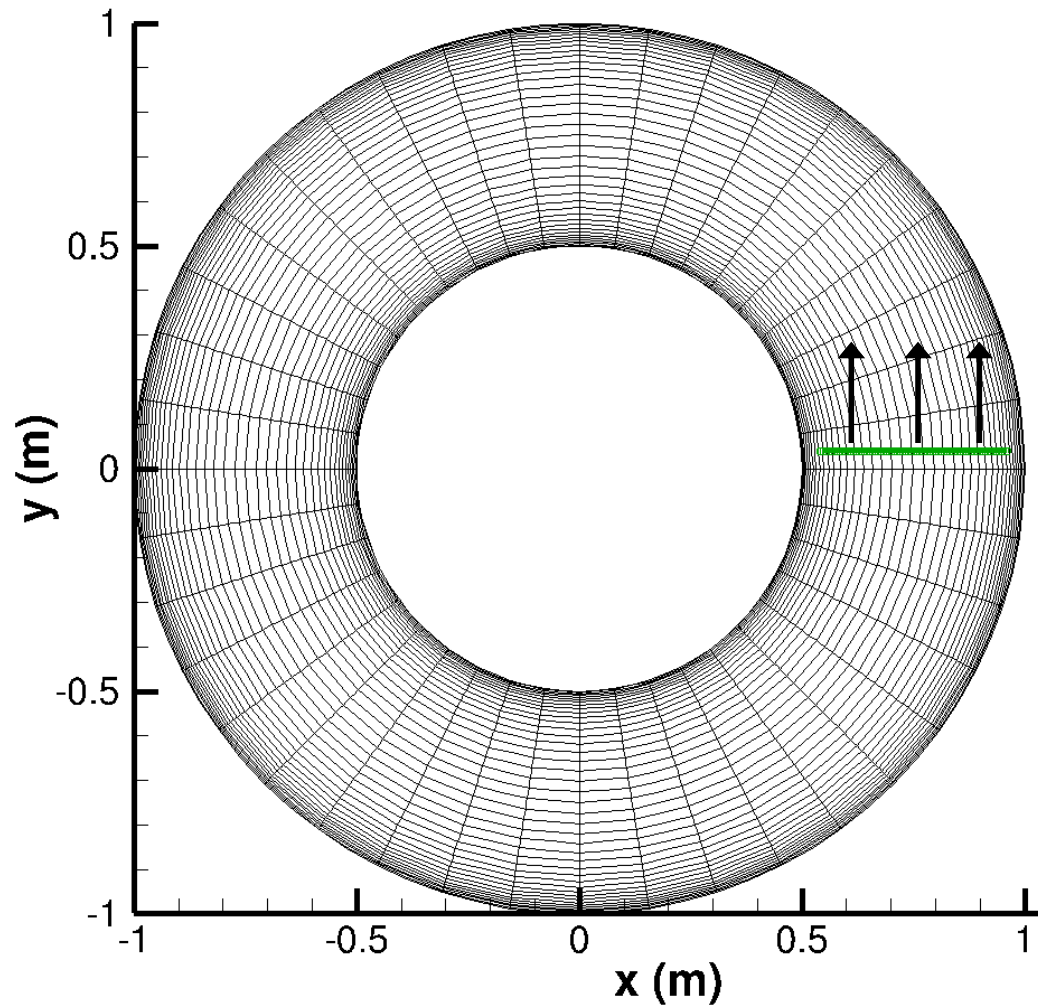


[1] Ching et al., *JCP*, 2020

[2] Ching and Ihme, *JCP*, submitted

Test case: particle advection in annular channel

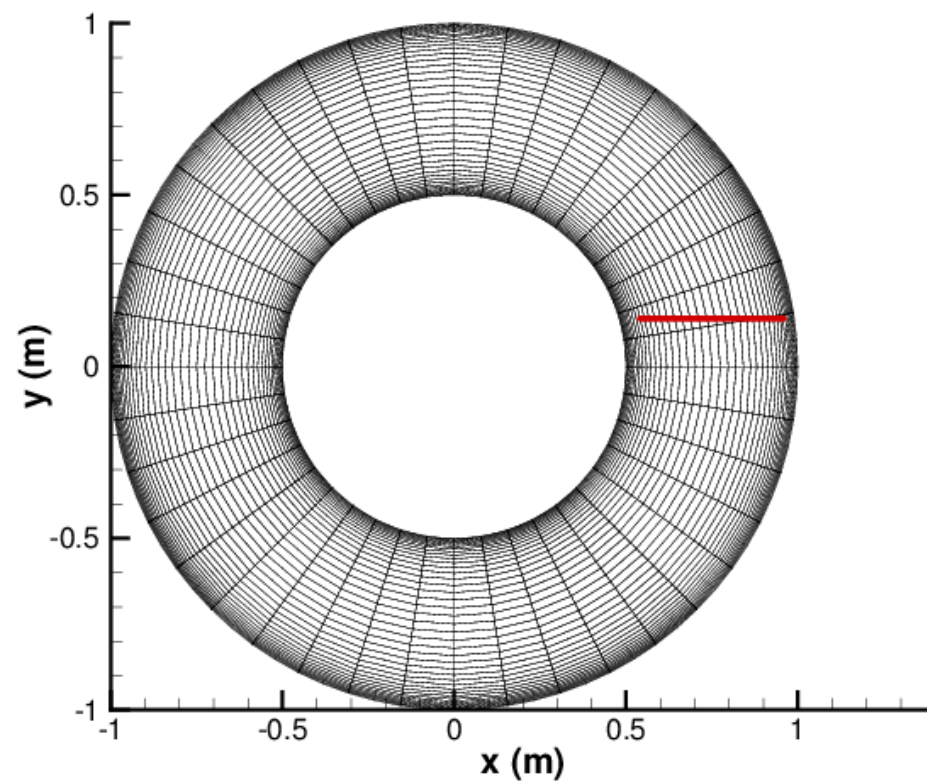
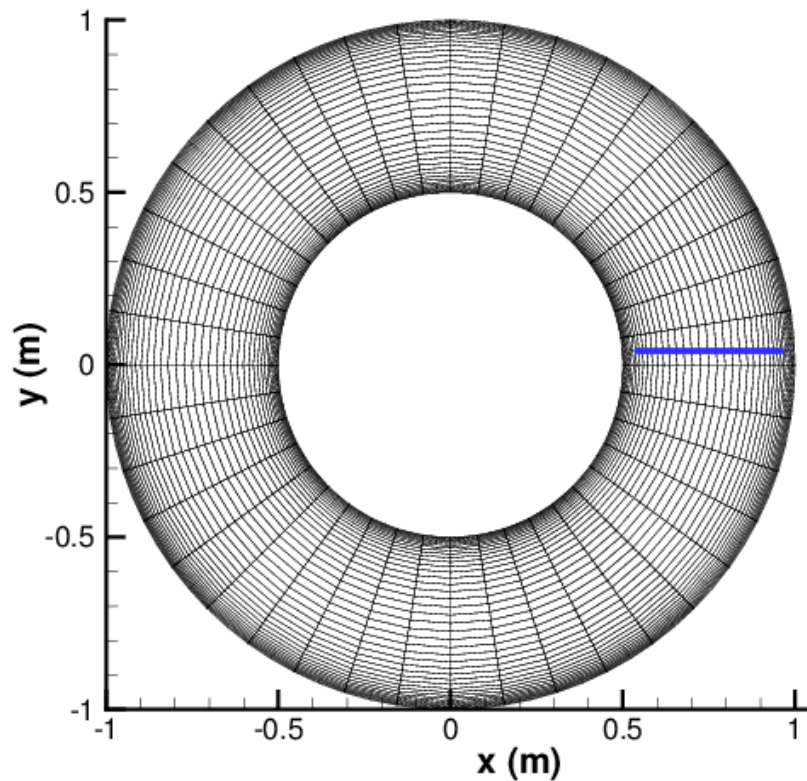
Compare straight-sided and curved elements



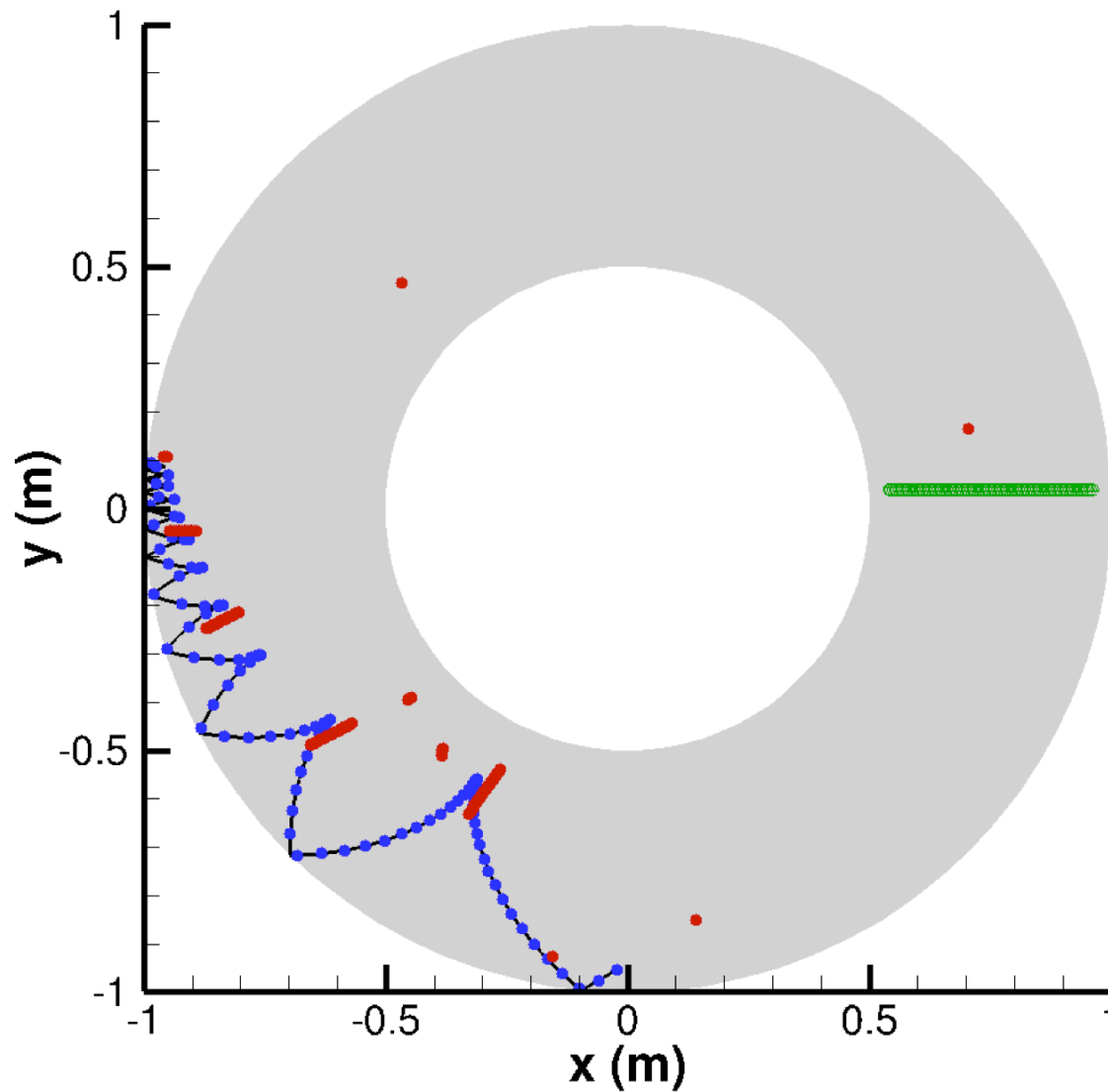
Test case: particle advection in annular channel

Blue: curved elements

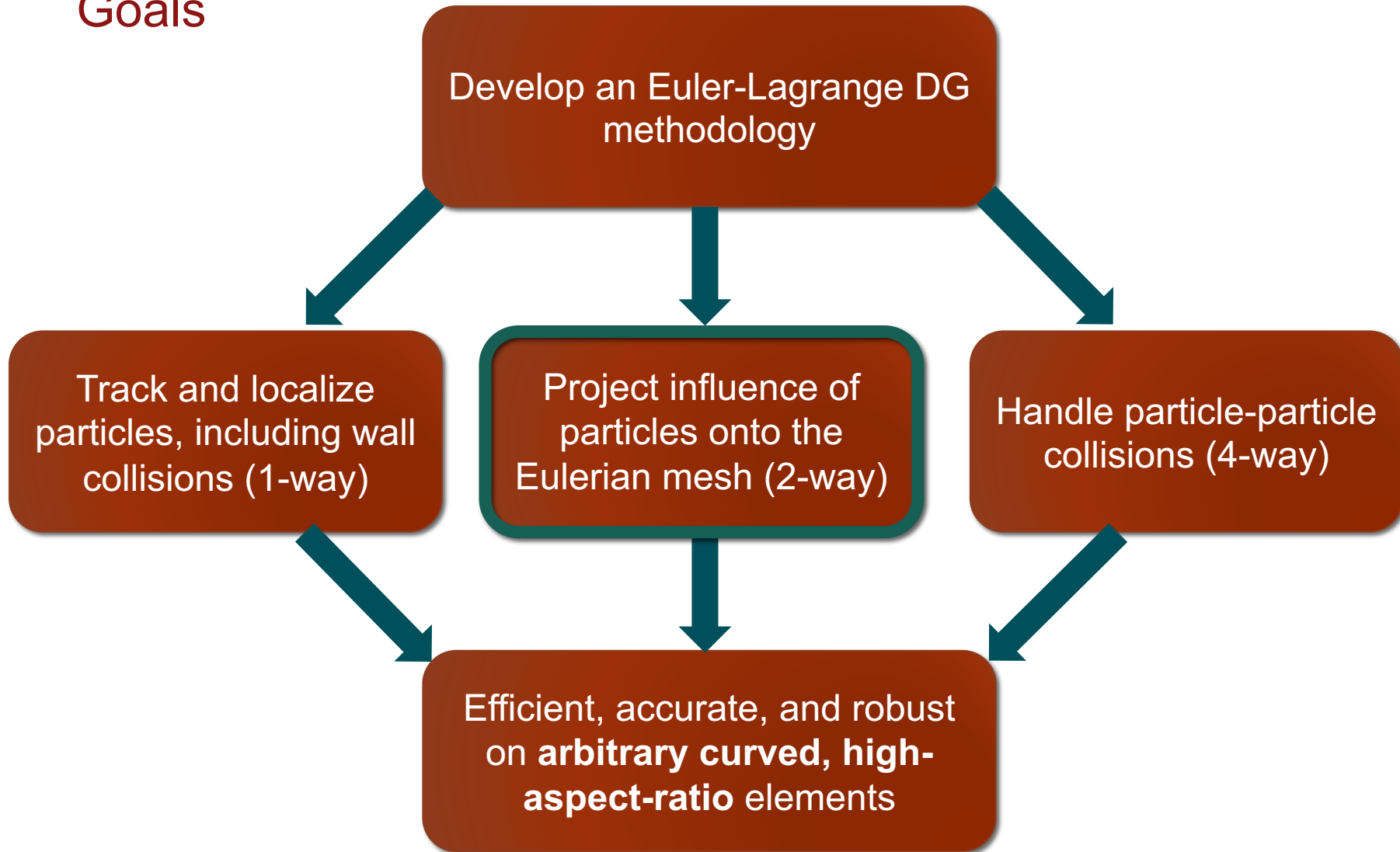
Red: straight-sided elements



Test case: particle advection in annular channel



Goals



Reverse coupling

Project the influence of the particles onto the Eulerian mesh

Source term $\mathbf{S} = [0, \mathbf{S}_m, S_e]^T$

$$\mathbf{S}_m = - \sum_{i=1}^{N_p} \mathbf{F}_i \chi(\mathbf{x}; \mathbf{x}_{d,i}),$$

Projection kernel $\chi(\mathbf{x}; \mathbf{x}_{d,i})$

- $\int_{\Omega} \chi d\Omega = 1$
- $\arg \max_{\mathbf{x} \in \Omega} \chi(\mathbf{x}; \mathbf{x}_d) = \mathbf{x}_d$

$$S_E = - \sum_{i=1}^{N_p} \left(Q_i + \mathbf{u}_{d,i} \cdot \mathbf{F}_i \right) \chi(\mathbf{x}; \mathbf{x}_{d,i})$$

Finite-volume context

- Box kernel¹
- Volume-weighted², distance-weighted³ kernel
- Isotropic Gaussian function⁴

[1] Crowe, *J Fluids Eng*, 1982

[2] Squires and Eaton, *Phys Fluid A*, 1990

[3] Elghobashi and Truesdell, *Phys Fluid A*, 1993

[4] Capecelatro and Desjardins, *J Comp Phys*, 2013

Reverse coupling

DG context

- Quadrature used to evaluate integrals
- Multiple DOFs per element, subcell resolution
- Modal basis functions
- More susceptible to numerical noise and instabilities
- Isotropic Gaussian function^{1,2,3}

Shifted delta function⁴ $\chi(\mathbf{x}; \mathbf{x}_{d,i}) = \delta(\mathbf{x} - \mathbf{x}_{d,i})$

$$\int_{\Omega_e} \phi_m \mathbf{S} d\Omega_e = \sum_{i=1}^{N_p^e} \phi_m(\mathbf{x}_{d,i}) [0, \mathbf{F}_i, Q_i + \mathbf{u}_{d,i} \cdot \mathbf{F}_i]^T$$

- Simple
- Not smooth \rightarrow numerical noise
- Inappropriate for larger particles

[1] Jacobs and Hesthaven, *J Comp Phys*, 2006

[2] Huang et al, *Comput Methods Appl Mech Eng*, 2019

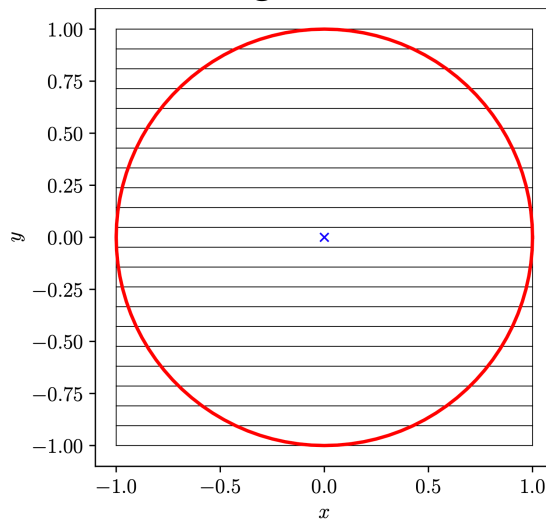
[3] Zwick and Balachandar, *Int J High Perform Comput Appl*, 2020

[4] Ching et al., *J Comp Phys*, 2020

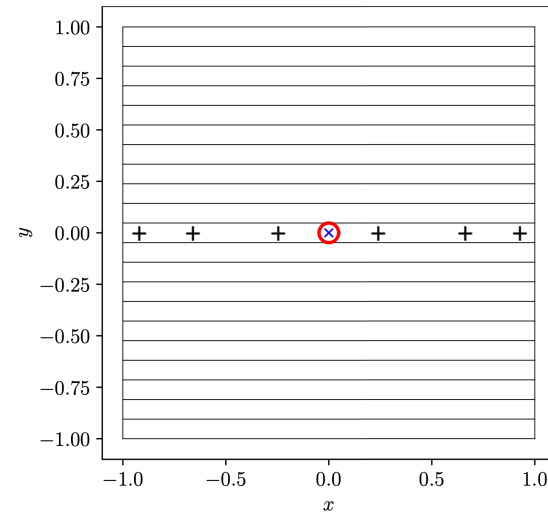
Smooth isotropic kernels: challenges

High-aspect-ratio elements

Large kernel

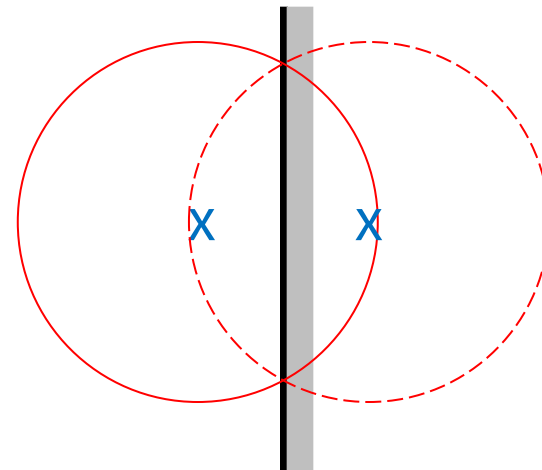


Small kernel



Curved elements

- To conserve **near-wall interphase transfer**, common to use mirror particles on straight-sided elements
- Not straightforward for arbitrary curved elements



Proposed methodology

1. Smooth anisotropic kernels

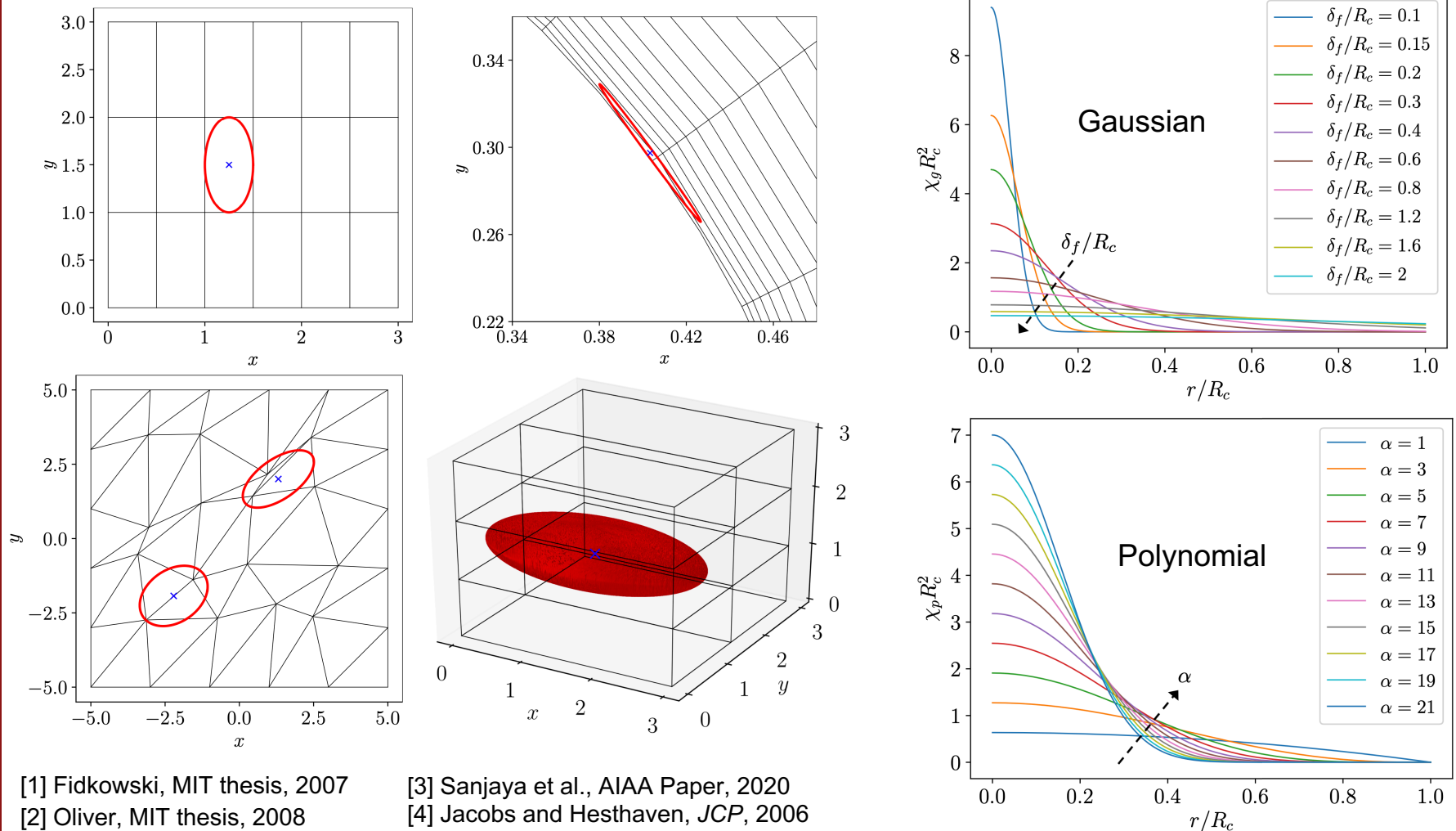
- Balance:
 - Accuracy
 - Mitigation of numerical noise and instabilities
 - Computational cost
- Elliptical/ellipsoidal (instead of circular/spherical) in 2D/3D

2. Efficient kernel rescaling near walls

- Avoid brute-force approach
- Employ high-order polynomials to approximate rescaling factor

Smooth anisotropic kernel

- Construct ellipse based on mesh-implied metric^{1,2,3}
- Polynomial-based kernel⁴ instead of Gaussian-type kernel



Near-wall treatment

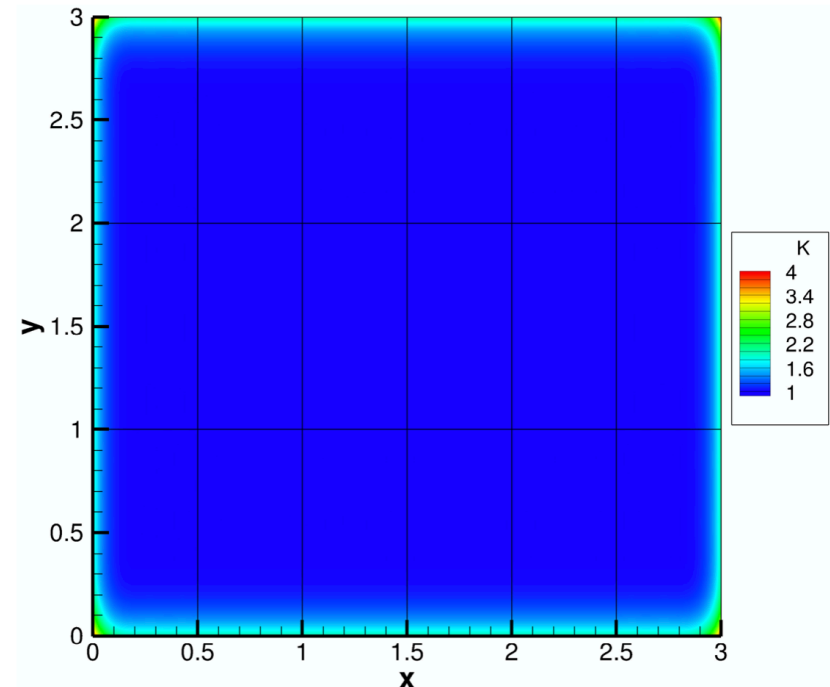
Brute-force approach: exact rescaling

$$K(\mathbf{x}_d) = \frac{1}{\sum_{e=1}^{N_v} \int_{\Omega_e} \chi_p(\mathbf{x}; \mathbf{x}_d) d\Omega}$$

Proposed approach: approximate rescaling

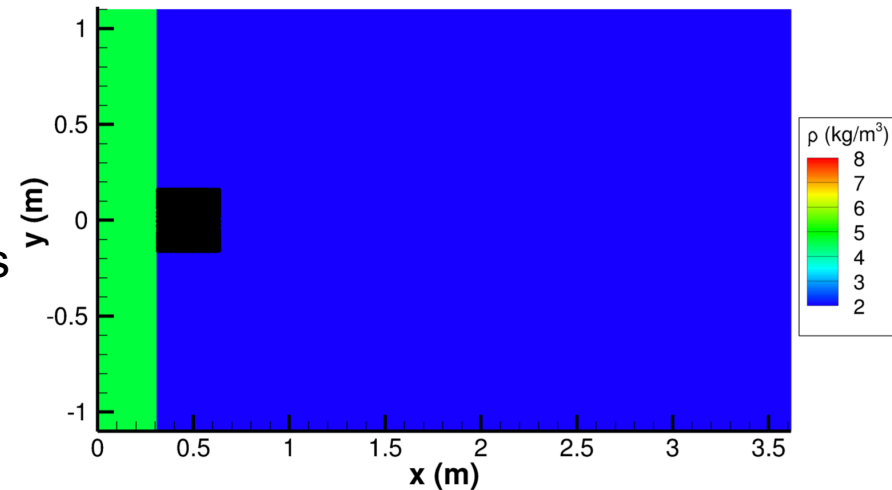
- Construct high-order polynomial approximation of K
 - Sample K at quadrature points
 - Project to high-order polynomials

$$K_h^e(\mathbf{x}_d) = \sum_{n=1}^{N_K} \tilde{K}_n^e \Phi_n(\mathbf{x}_d)$$



Test case: shock interaction with particle cloud

- Bronze particle cloud
- Volume fraction = 4%
- 500,000 2D quadrilateral elements
- $p = 3$ polynomials

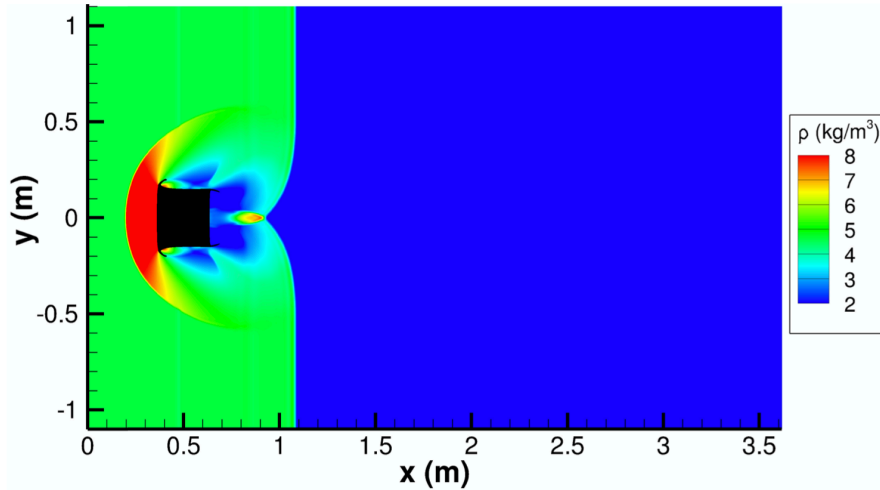


ρ_d (kg/m ³)	D (μm)	c_d (J/kg/K)	ρ_c (kg/m ³)	u_c (m/s)	P_c (bar)
8900	100	435	1.2	0	1

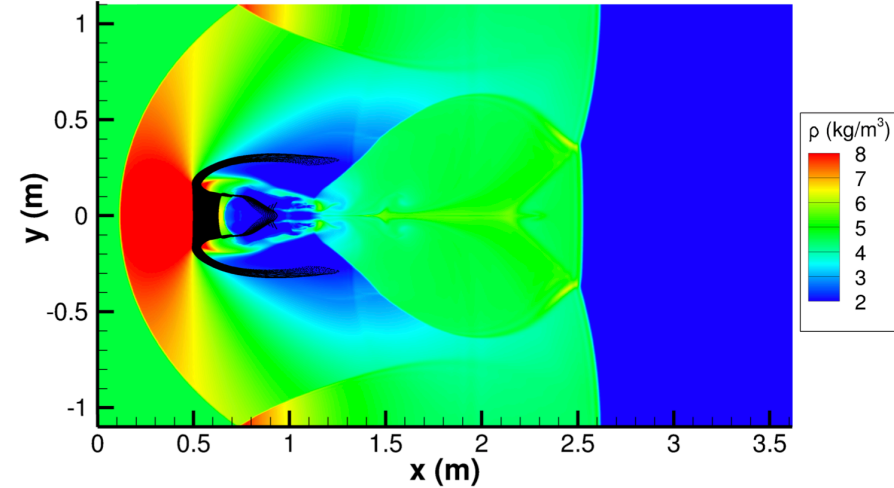
- Difference between delta shape function and smooth anisotropic kernel
- Good agreement with numerical results¹

Test case: shock interaction with particle cloud

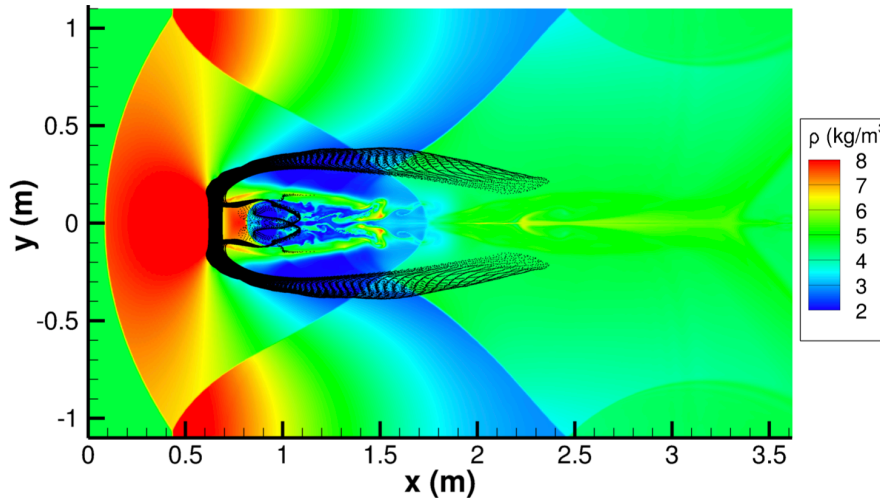
$t = 0.75$ ms



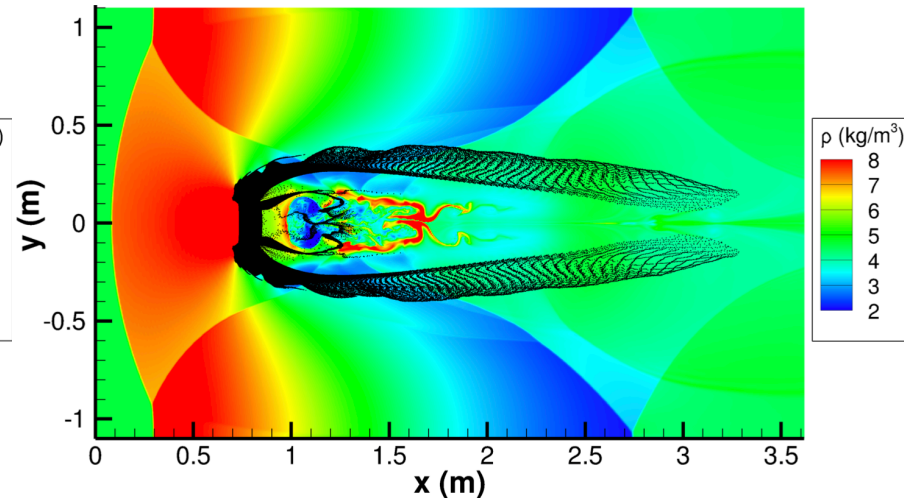
$t = 2.25$ ms



$t = 3.75$ ms

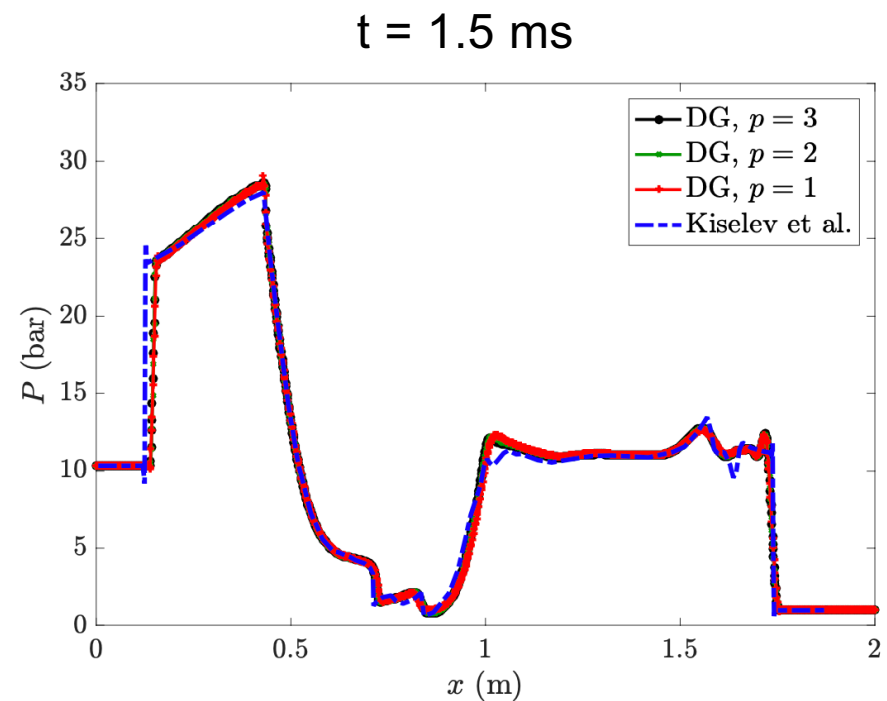
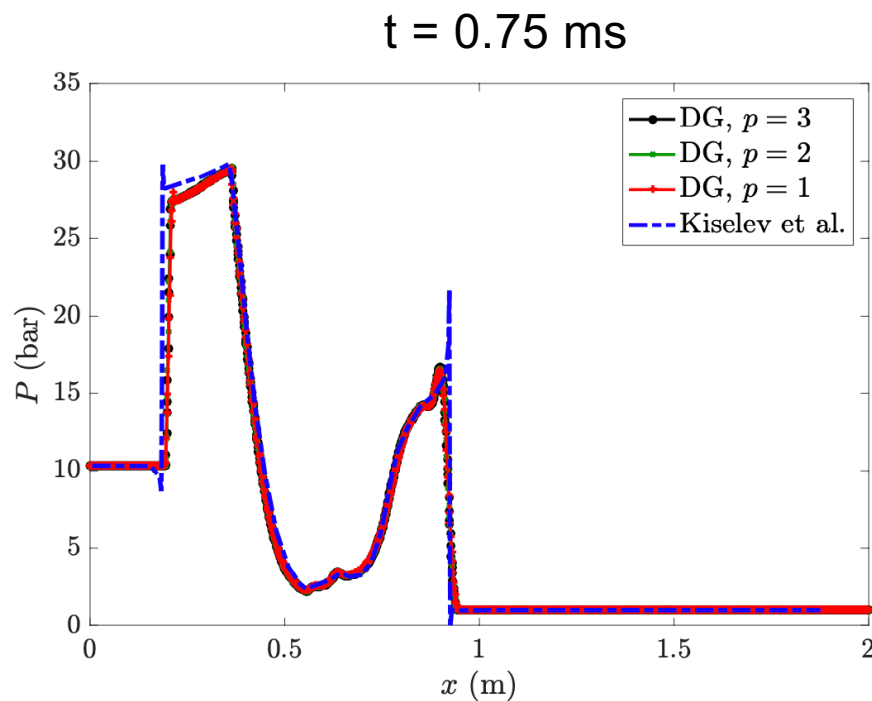


$t = 5.0$ ms



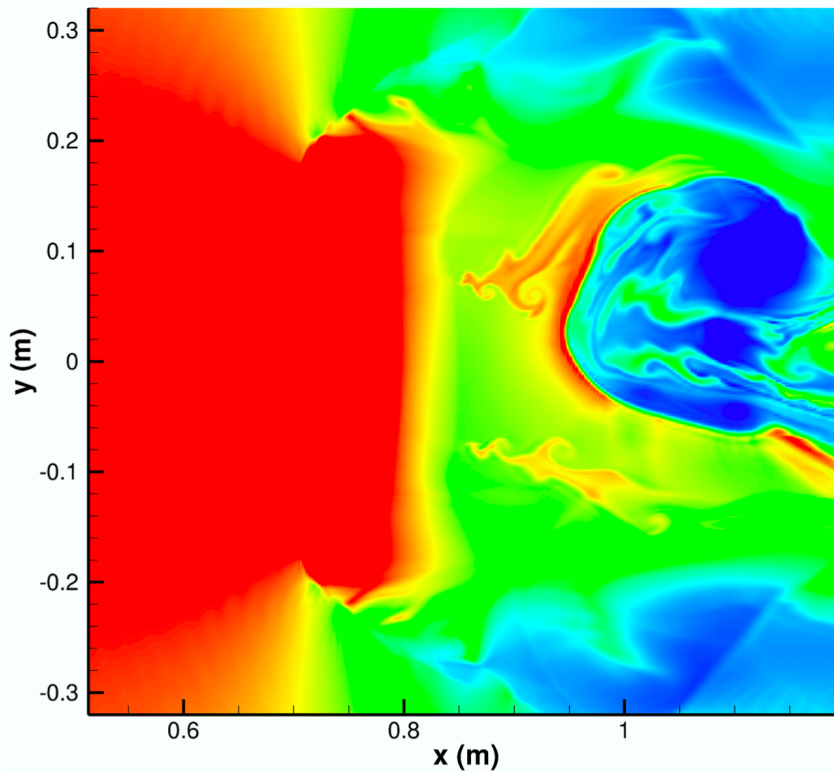
Test case: shock interaction with particle cloud

Comparison with numerical results

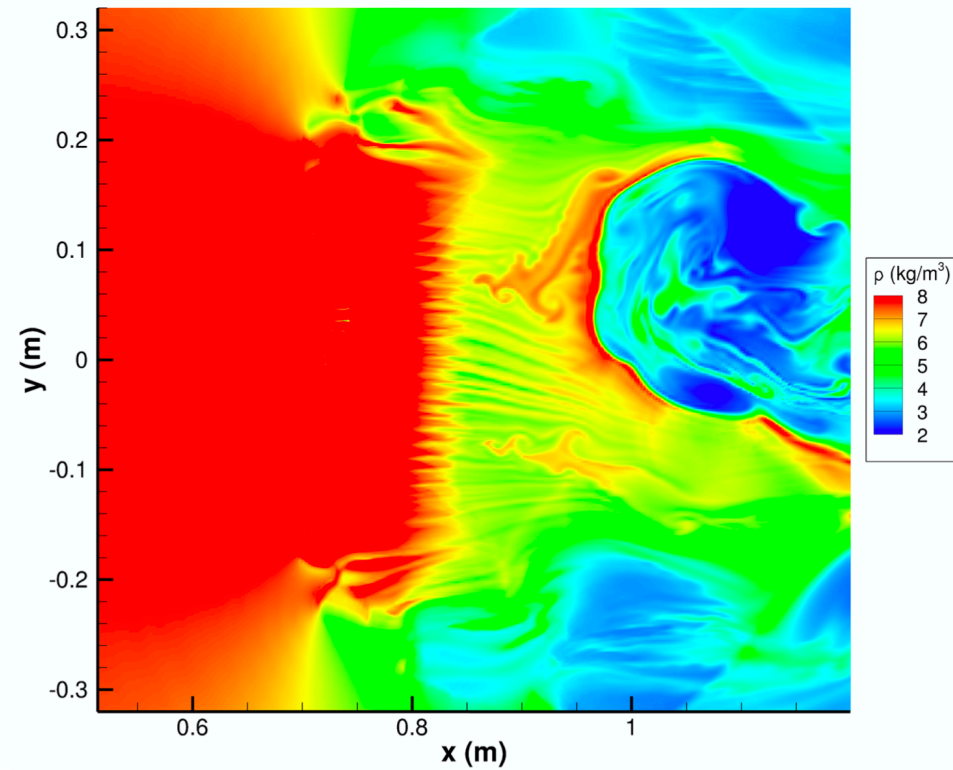


Test case: shock interaction with particle cloud

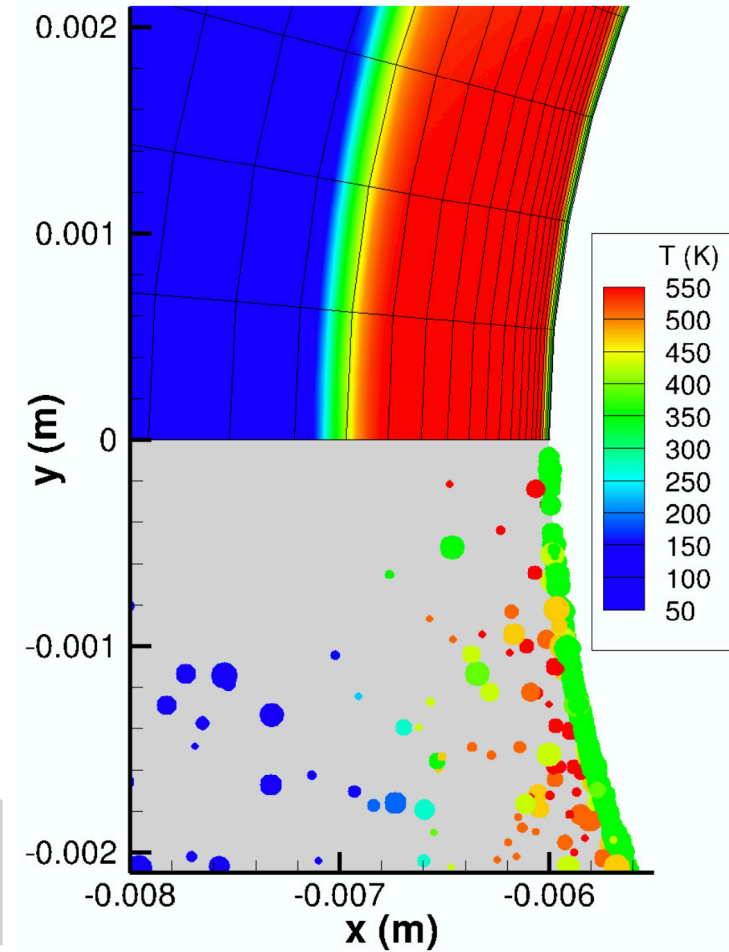
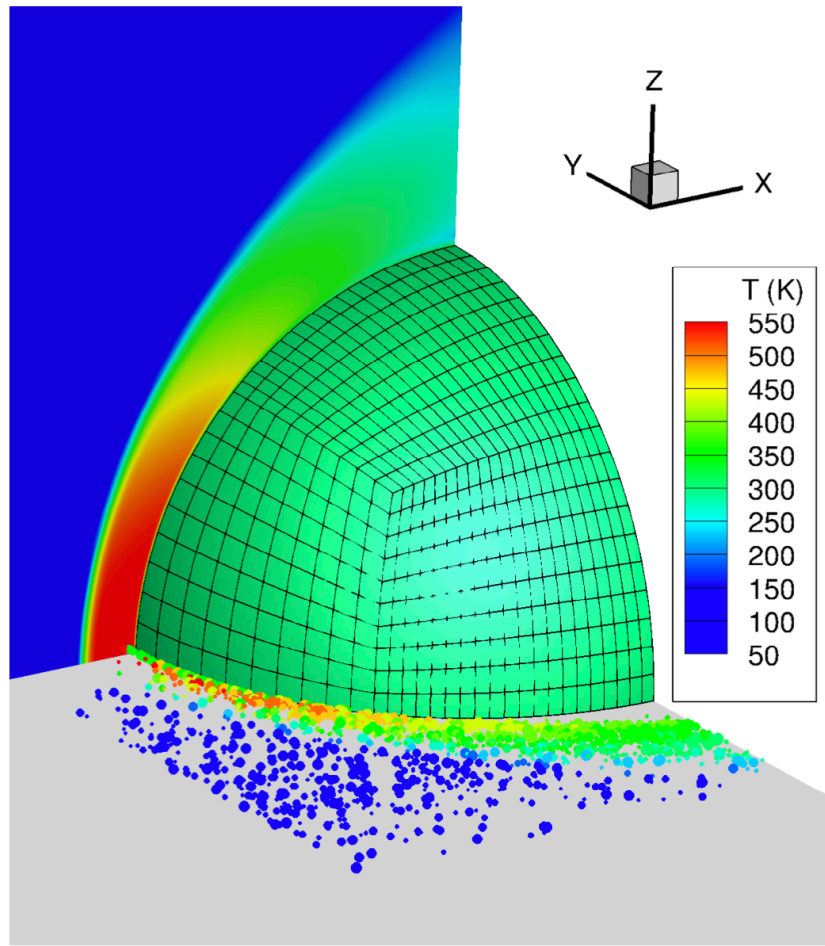
Smooth anisotropic kernel



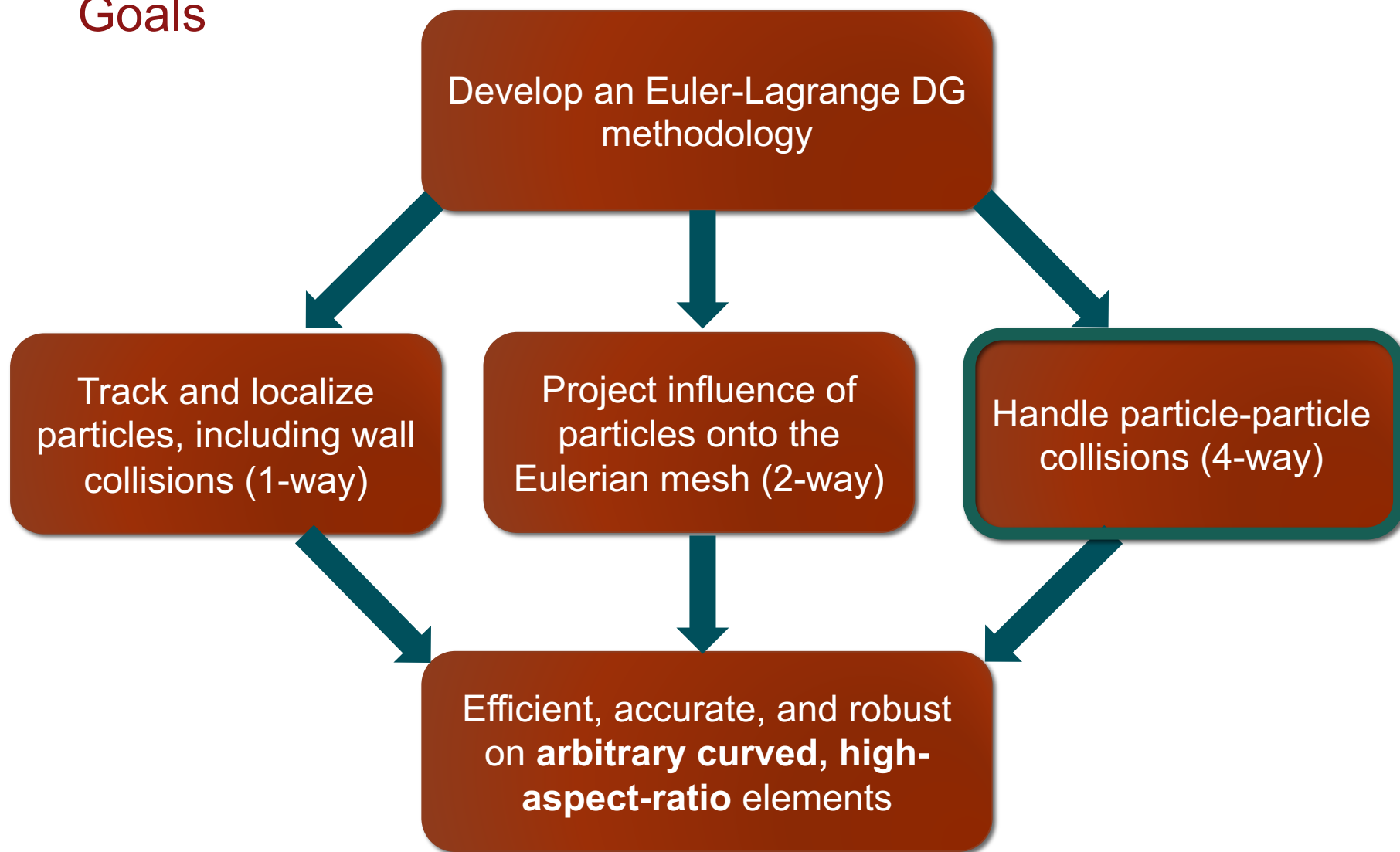
Delta functions



Hypersonic dusty flows over blunt bodies

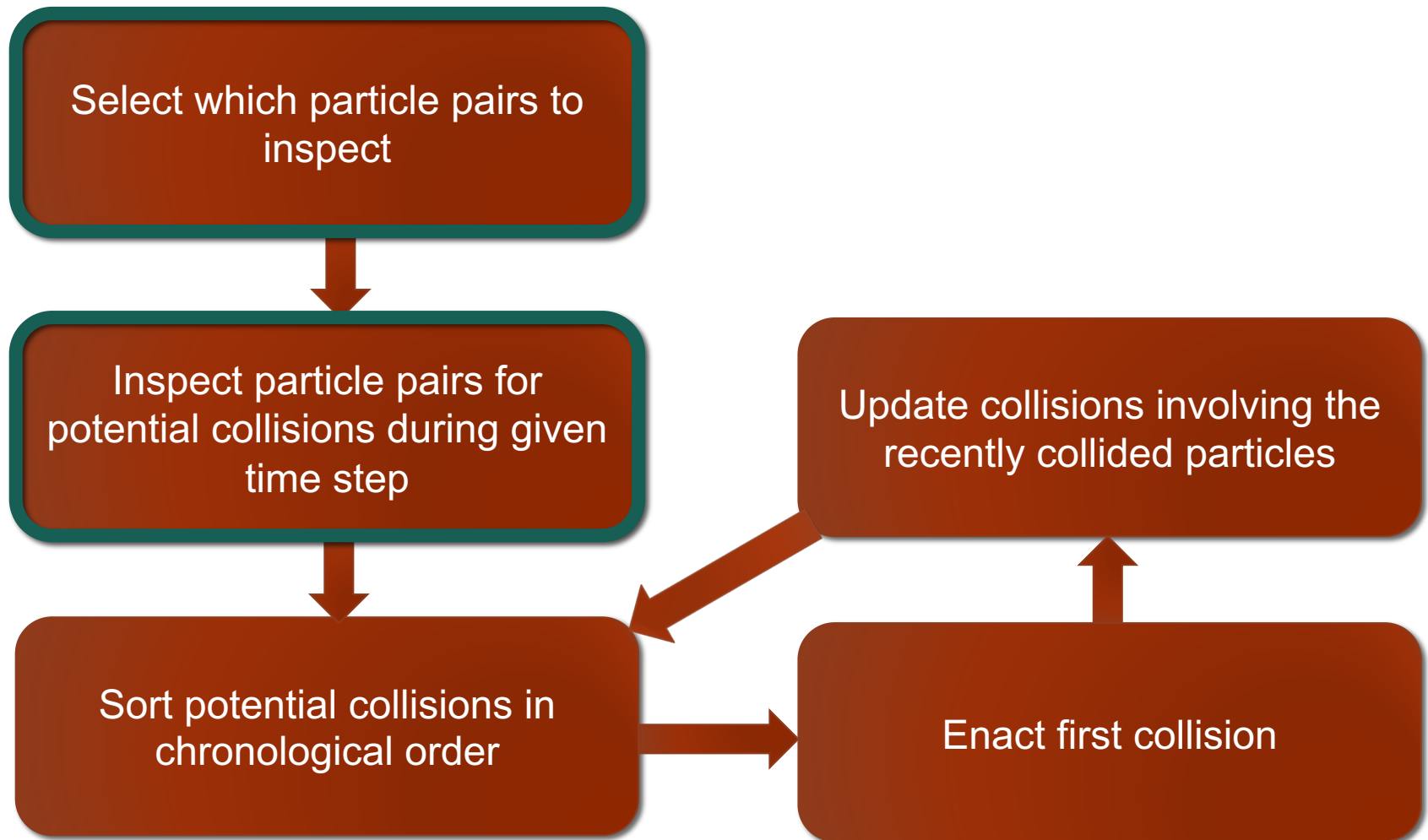


Goals



Particle-particle collisions

Hard-sphere model

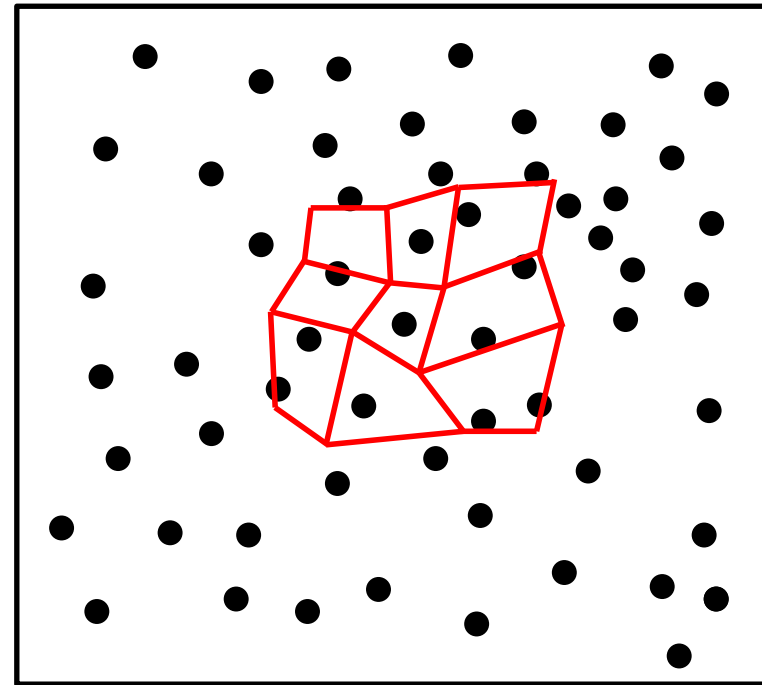


Selecting which particle pairs to inspect

- Brute-force approach
 - All possible particle pairs in domain
- Element neighbor lists
 - Inspect particles in node-sharing elements

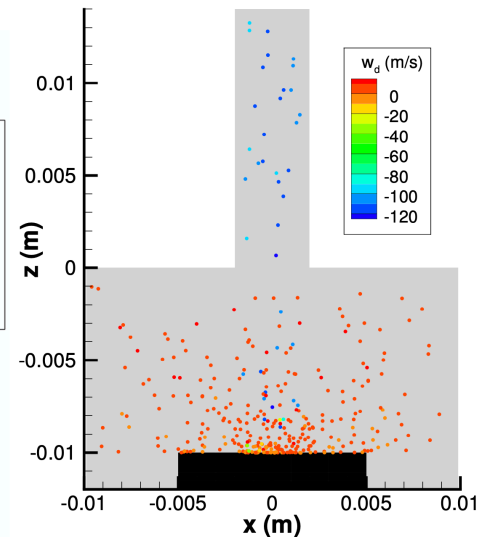
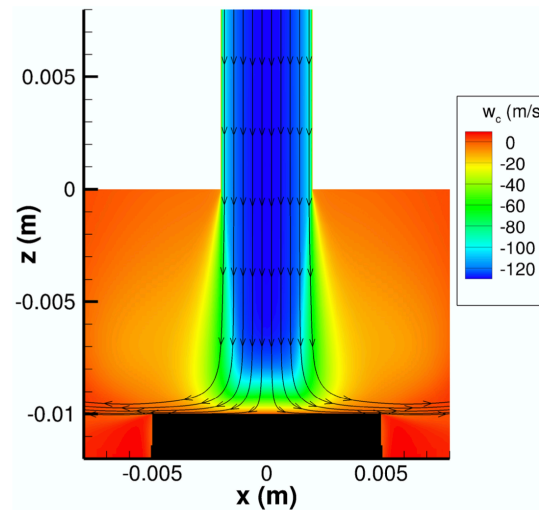
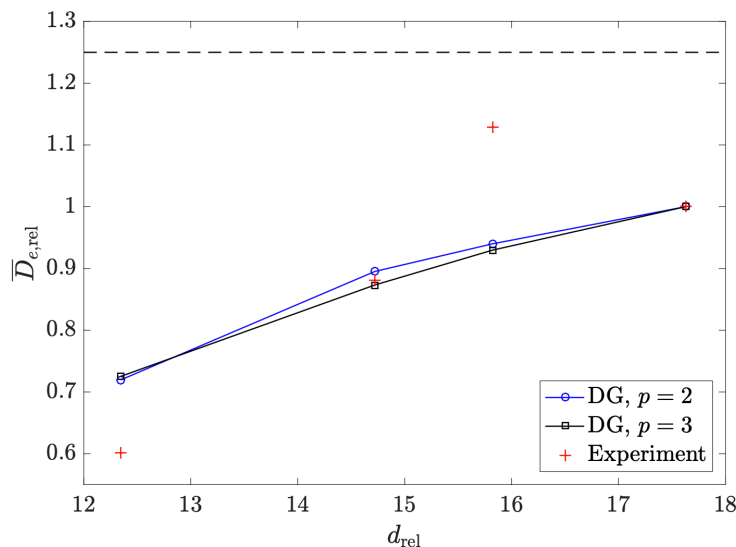
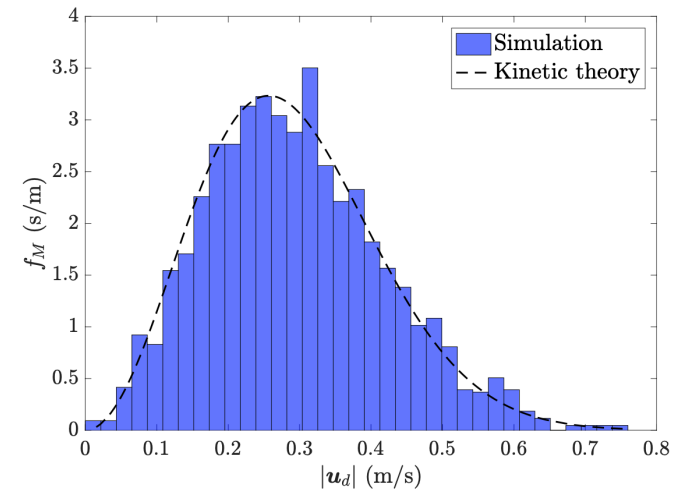
Truncate element neighbor lists

- Exploit the geometric mapping to further localize pairings
- Compatible with unstructured curved elements
- Simple and fast truncation process



Test cases

- 2D/3D kinetic theory example
 - Recover equilibrium properties
 - Significant speedup
- 3D sandblasting example
 - Impingement of particle stream on a flat plate
 - Effect of particle mass flux on erosion
 - Good agreement with experiments¹



Summary

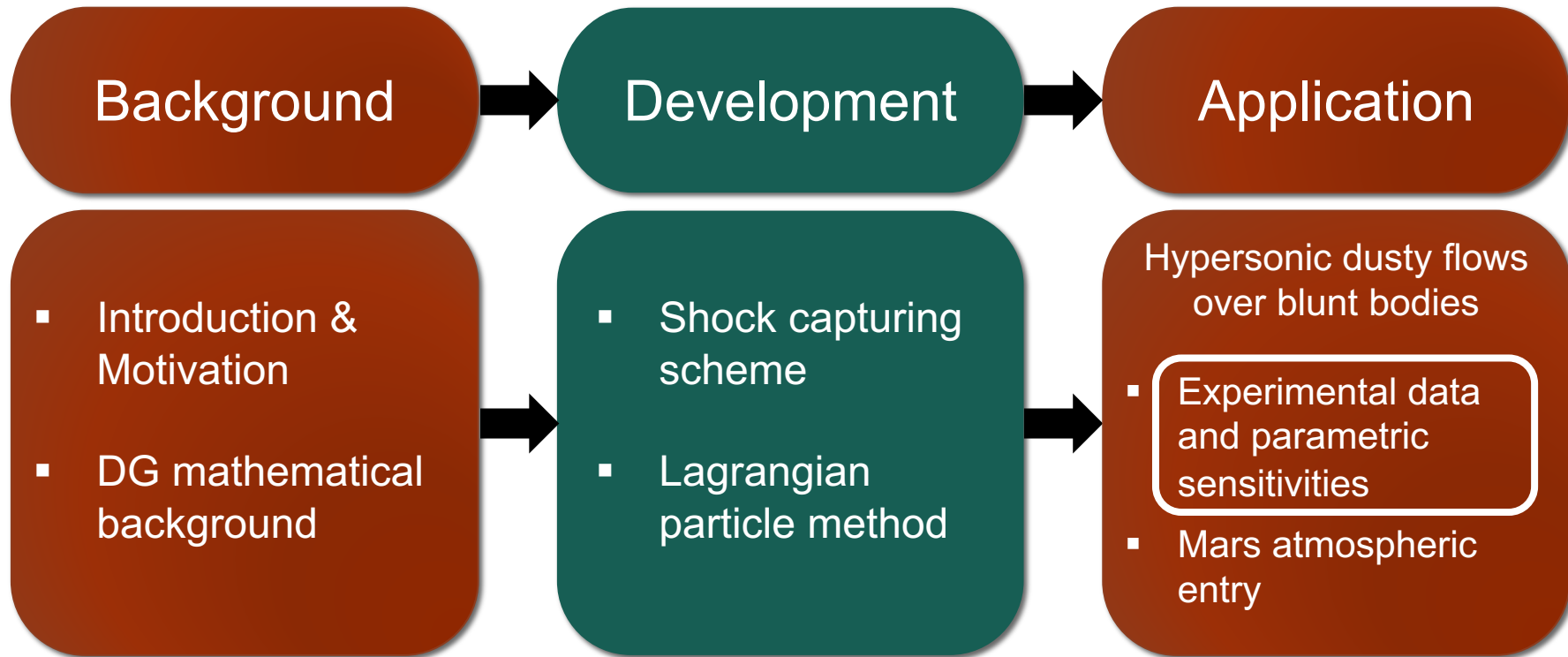
Developed an Euler-Lagrange DG solver

- Simple and robust shock capturing
 - Accurate surface heating predictions in hypersonic flows
- Lagrange particle methods for 1-way, 2-way, 4-way coupling
 - Compatible with arbitrary curved, high-aspect-ratio elements

Why DG?

- DG can considerably simplify meshing of complex geometries
- Curved elements can significantly improve predictions of particle trajectories

Outline



Hypersonic dusty flows over blunt bodies

Relevant to Mars atmospheric entry

Adverse effects¹:

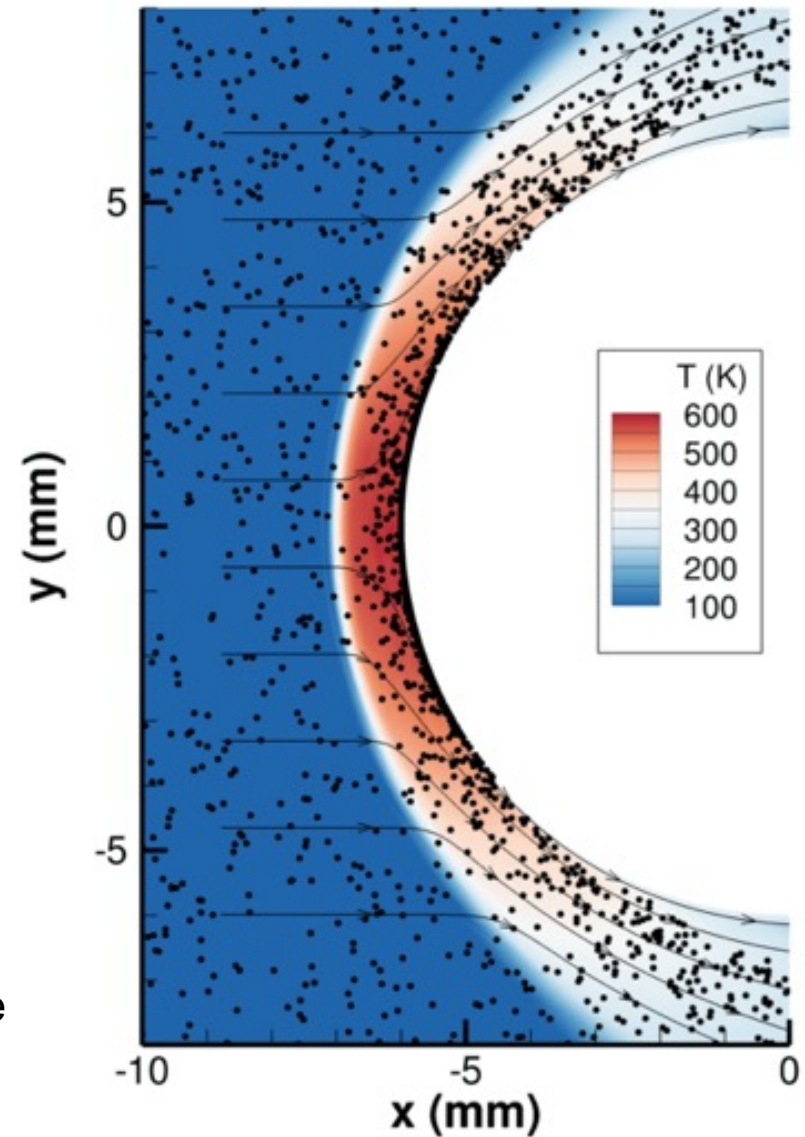
- Surface erosion
 - Particle-wall collisions
- Surface heat flux augmentation
 - Particle-wall collisions
 - Two-way coupling

Limited high-quality experimental data

- Ambiguity in how to reliably model these flows

Goals:

- Simulate the hypersonic dusty flow experiment by Vasilevskii et al.²
- Evaluate solution sensitivities to physical modeling of the particle phase



[1] Montois et al., International Planetary Probe Workshop, 2007

[2] Vasilevskii et al., *J Eng Phys Thermophys*, 2001.

Experiments by Vasilevskii et al.

Hypersonic dusty flow over a sphere

Gas	Ma_∞	$P_{t,\infty}$ (bar)	$T_{t,\infty}$ (K)	R_s (m)
N ₂	6.1	17.5	570	0.006

Dust material	ρ_d (kg/m ³)	\overline{D} (m)	β
SiO ₂	2264	0.19×10^{-6}	0.03

Measured ratio of dusty-gas to pure-gas heat flux at stagnation point

Solver details

- $p = 2$
- 60,000 elements
- Implicit-explicit time stepping
 - BDF3-AB3
- Two-way coupling

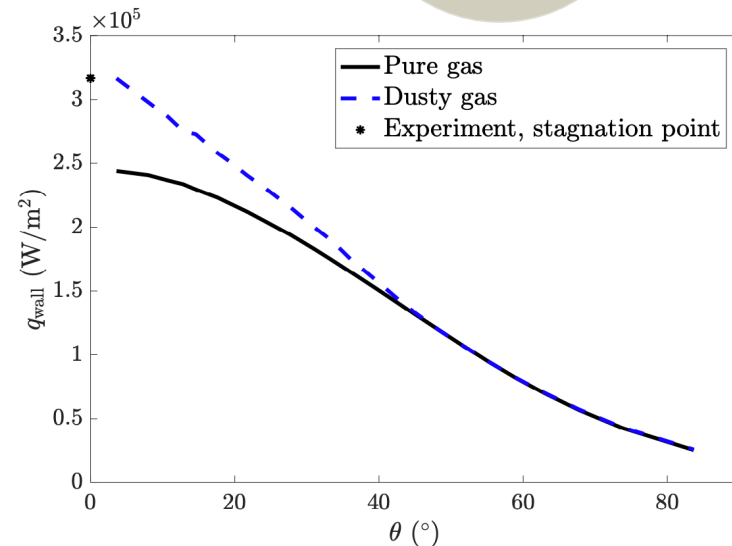
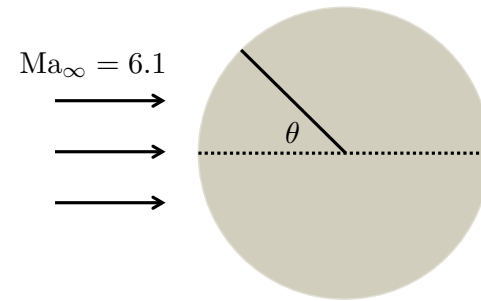
Baseline physical model

Simplest model that gives good agreement with experiment

- Henderson drag correlation¹
- Fox Nusselt number correlation²
- Thermophoretic force³

$$m_d \frac{d\mathbf{u}_d}{dt} = \mathbf{F} = \mathbf{F}_{qs} + \mathbf{F}_{th}$$

$$m_d c_d \frac{dT_d}{dt} = Q = Q_{qs}$$



Sensitivity study to be conducted with respect to this baseline model

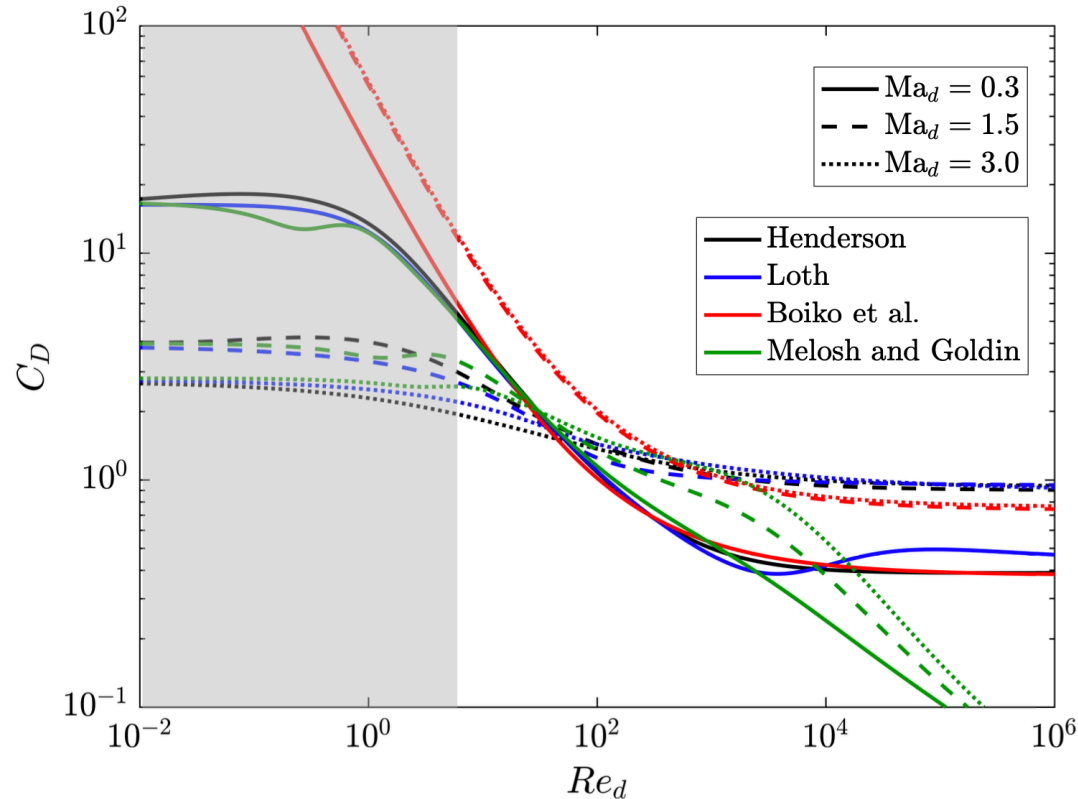
[1] Henderson, *AIAA J*, 1976

[2] Fox et al., 11th International Shock Tubes and Waves Symposium, 1978

[3] Loth, *AIAA J*, 2008

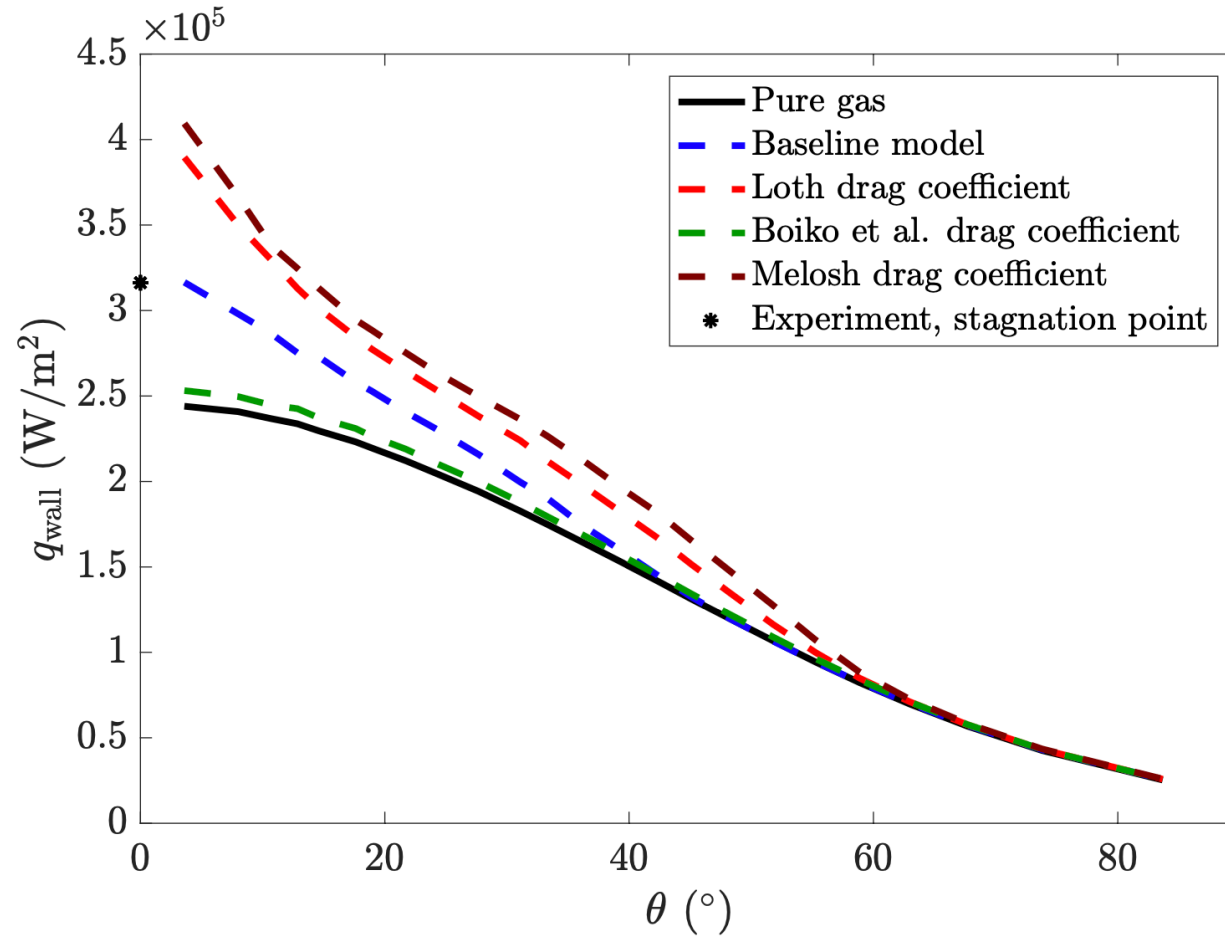
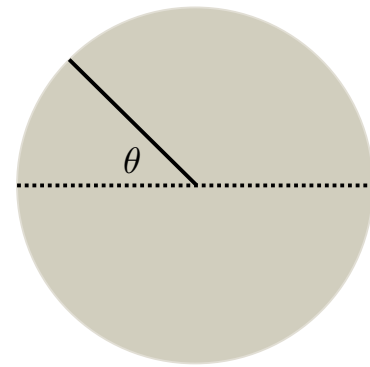
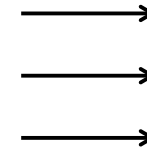
Drag correlations

- Henderson (*AIAA J*, 1976)
- Loth (*AIAA J*, 2008)
- Boiko et al. (*Shock Waves*, 1997)
- Melosh and Goldin (*Lunar and Planetary Science*, 2008)



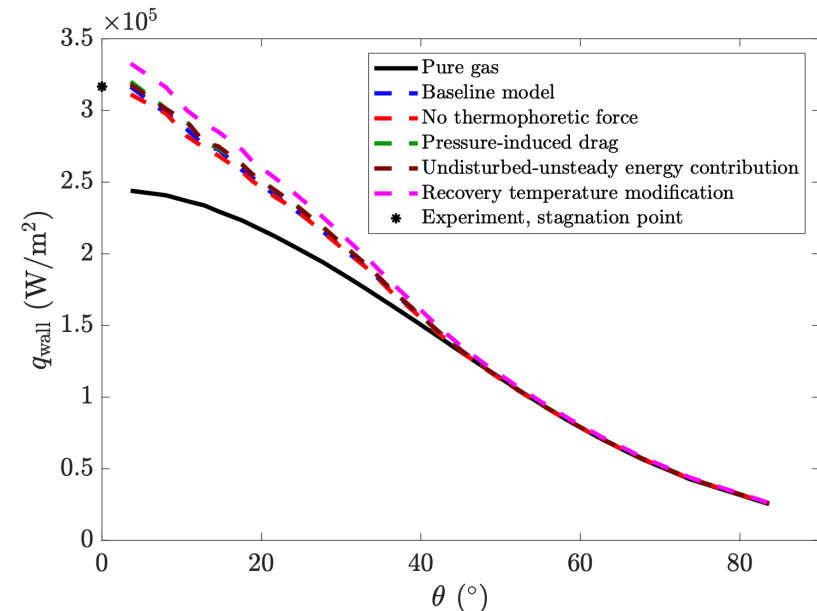
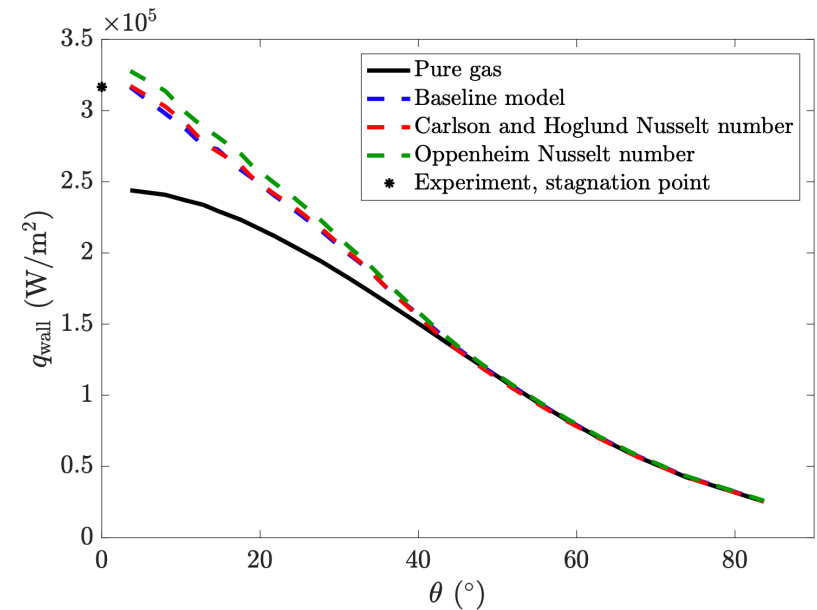
Sensitivity to drag correlation

$$\text{Ma}_\infty = 6.1$$



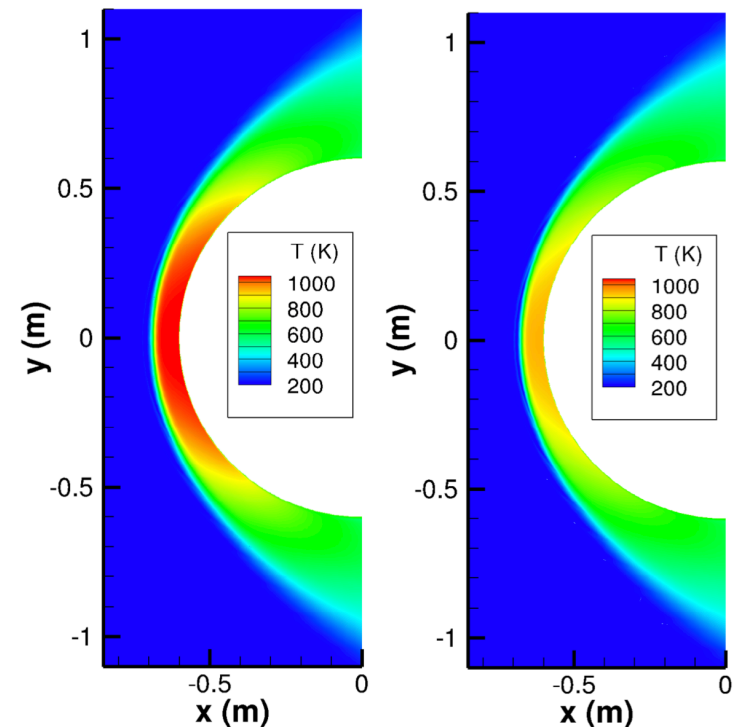
Additional sensitivities

- Nusselt number correlations
 - Small sensitivity
- Momentum and energy contributions
 - Quasi-steady drag and heating most important



Additional sensitivities

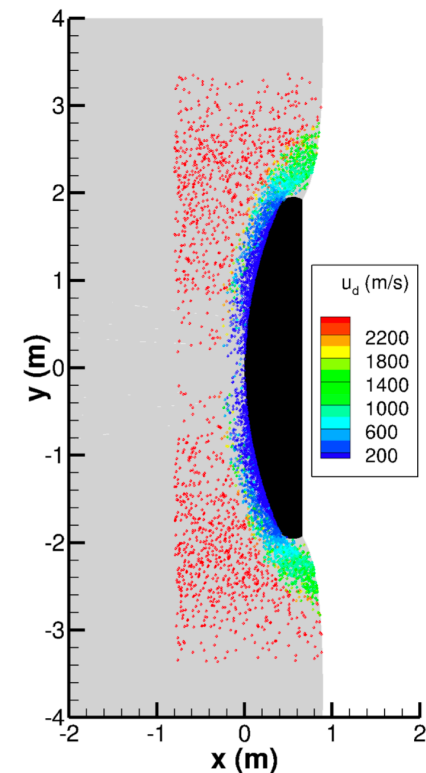
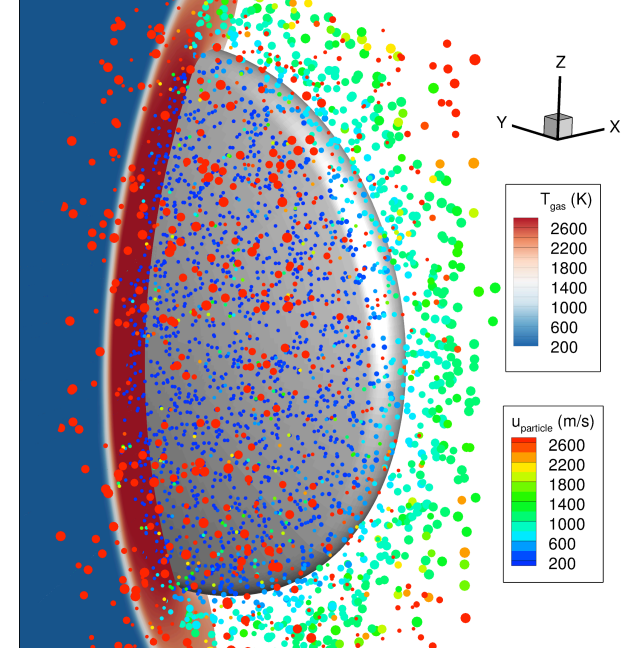
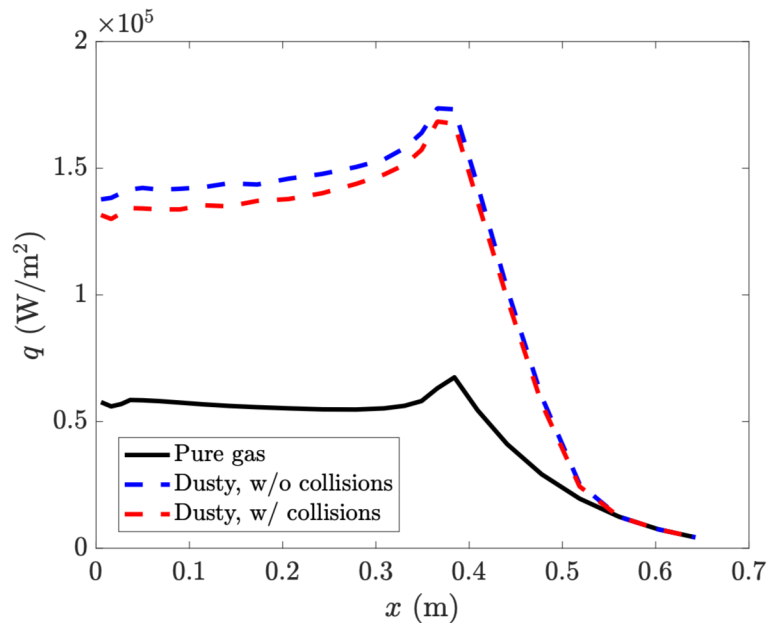
- Nusselt number correlations
 - Small sensitivity
- Momentum and energy contributions
 - Quasi-steady drag and heating most important
- Gas model
 - Shock standoff distance and shock layer quantities can have significant influence



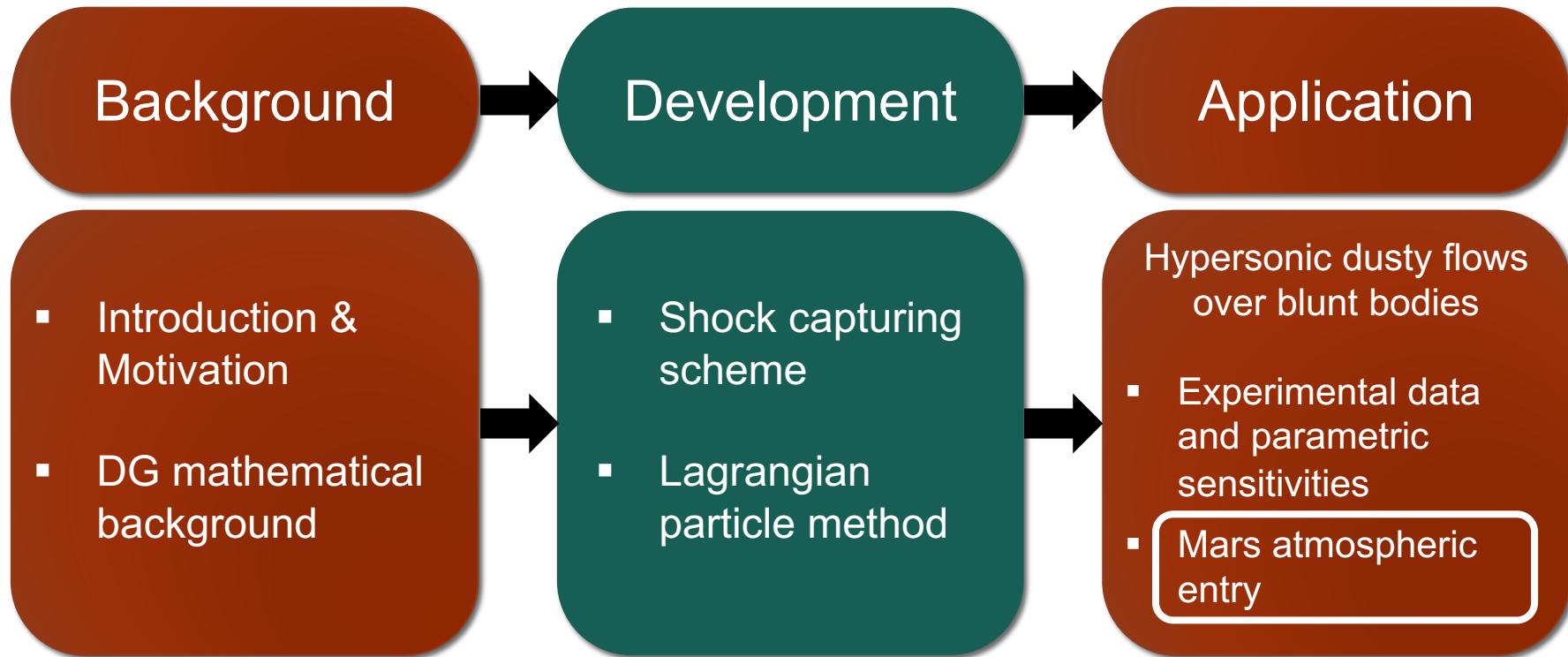
Additional sensitivities

Particle-particle collisions

- Hypersonic dusty flow over a capsule forebody
- Can attenuate particle-wall collisional energy flux
- Only important at very high dust loading



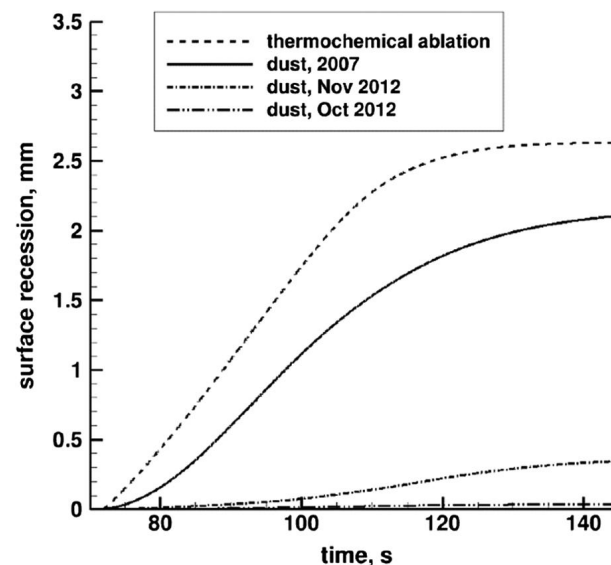
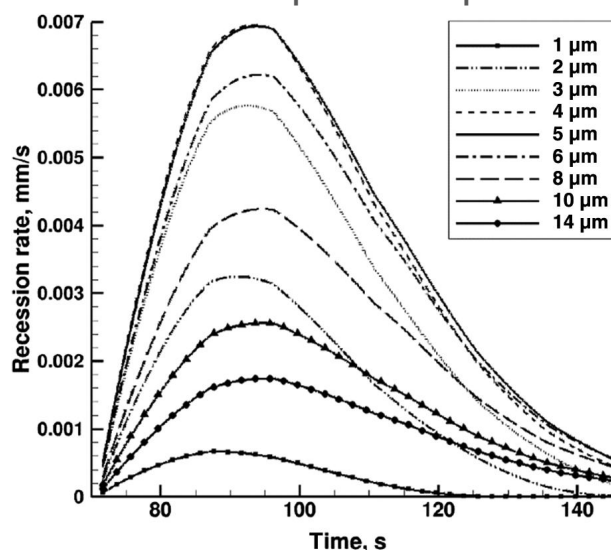
Outline



Background

Build on:

- Previous section
- Palmer et al., *Journal of Spacecraft and Rockets*, 2020
 - Computed heat shield recession at the stagnation point during the trajectory of the ExoMars Schiaparelli capsule



Goals:

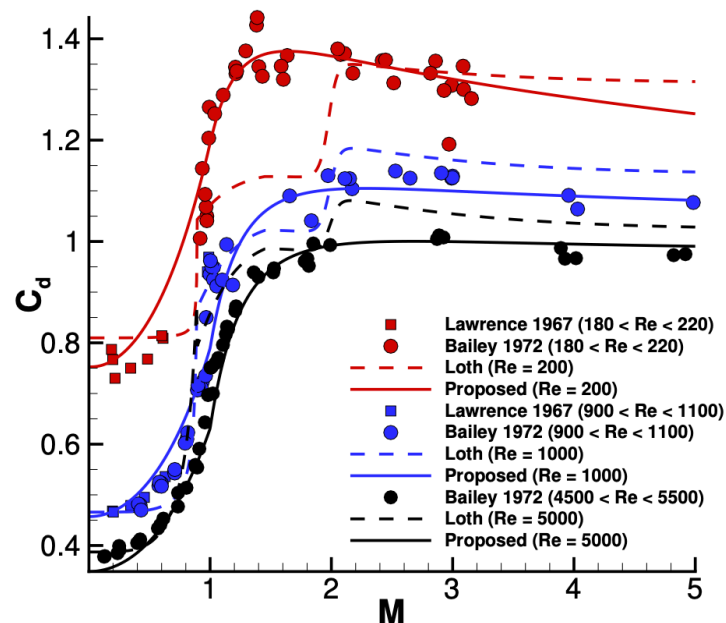
- Further improve understanding of relevant physical processes
- Assess effects of particle size distribution, angle of attack, and two-way coupling on heating augmentation and erosion
- Employ recently developed physics-based drag correlation¹
- Investigate particle trajectories in aft region

Drag correlation

Recently developed physics-based drag correlation¹ that incorporates:

- Low-speed hydrodynamics
- Shockwave physics
- Rarefied gas dynamics

Better agreement with experimental data than popular correlations^{2,3}



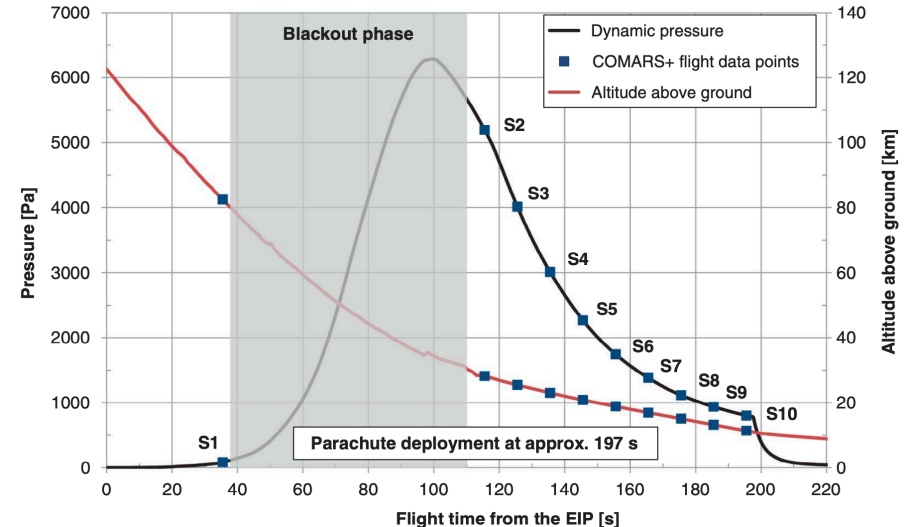
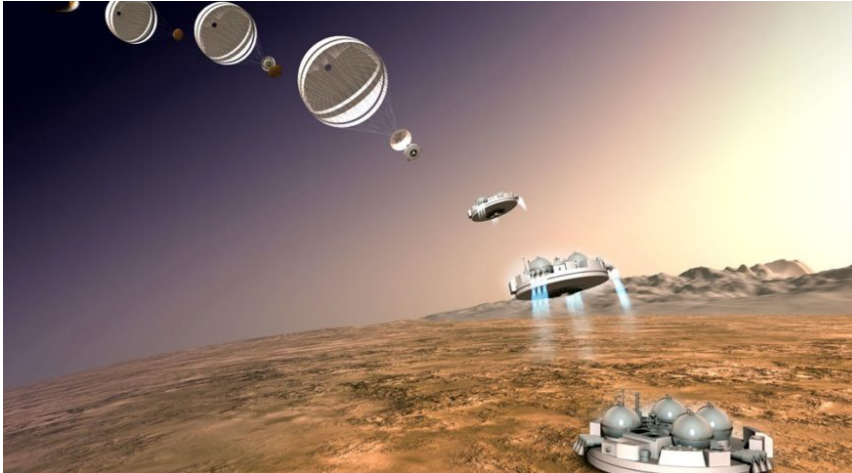
[1] Singh et al., <https://arxiv.org/abs/2012.04813>, 2020

[3] Loth, AIAA J, 2008

[2] Henderson, AIAA J, 1976

ExoMars Schiaparelli trajectory

- 1st mission of the European Space Agency's ExoMars program
- **Goals:**
 - Search for evidence of methane and other gases that could signify biological or geological processes
 - Test key technologies to support future missions



ExoMars Schiaparelli trajectory

S3 trajectory point¹

Gas	Mach	u (m/s)	P (Pa)	T (K)
CO ₂	8.97	2014	74.1	195

Two angles of attack: 0 and 20 deg

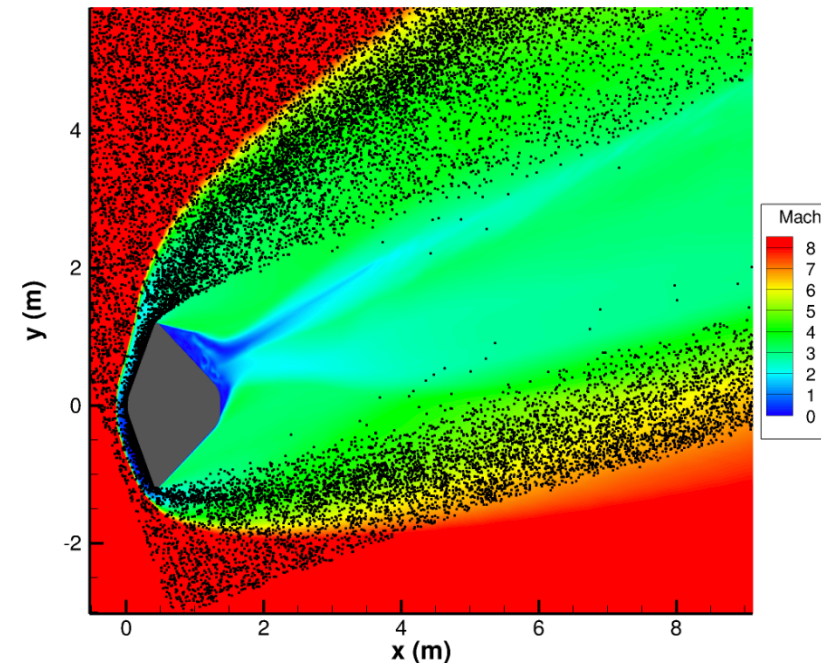
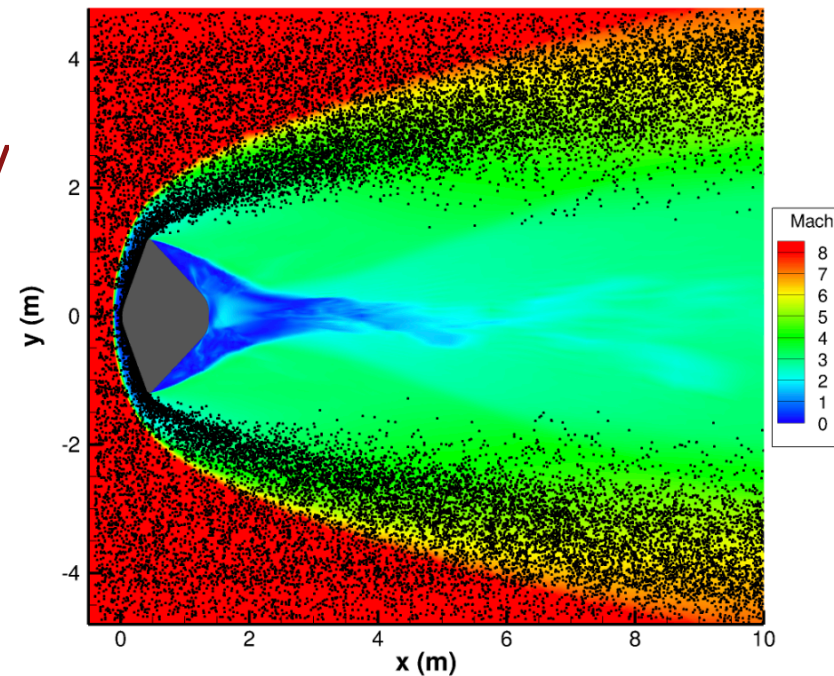
SiO₂ particles

Particle-wall impacts

- Heating augmentation by collisional energy transfer
- Surface recession²

~300,000 3D hexahedral elements

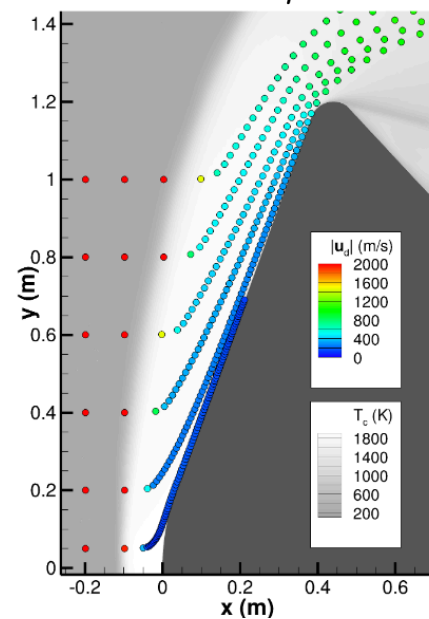
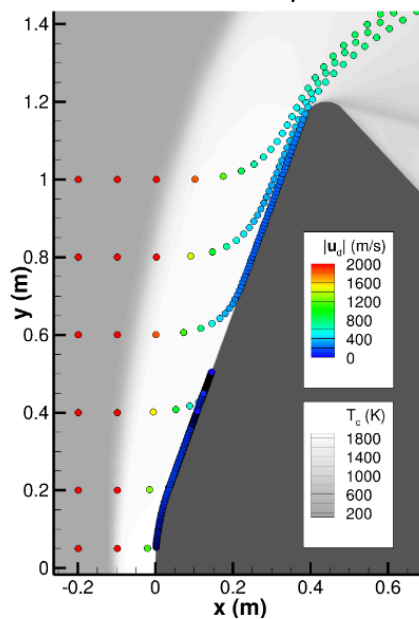
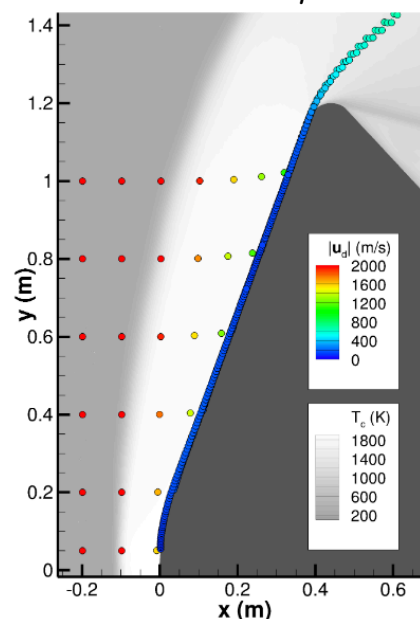
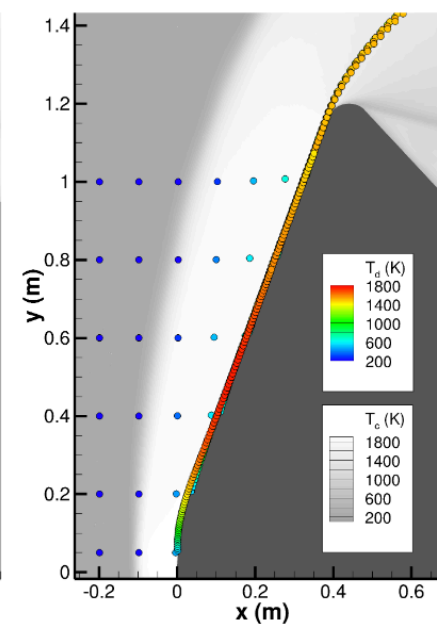
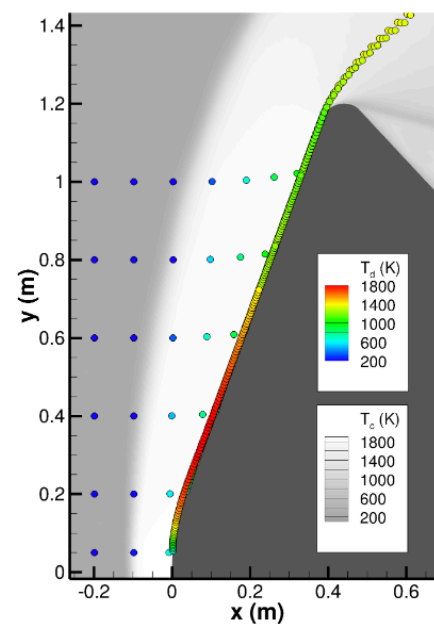
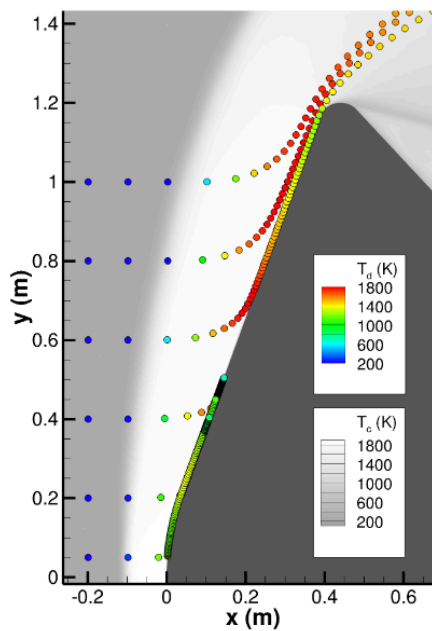
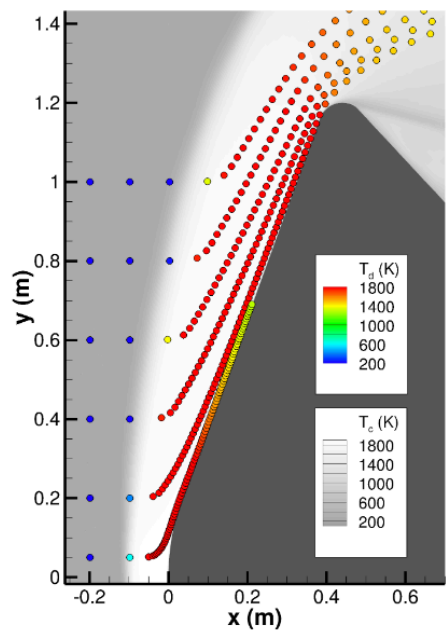
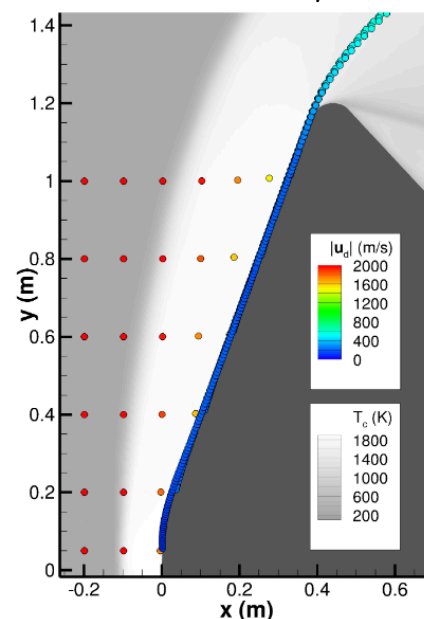
Third-order-accurate ($p = 2$)
polynomials



[1] Gulhan et al., *J Spacecraft Rockets*, 2019

[2] Palmer et al., *J Spacecraft Rockets*, 2020

Sample particle trajectories

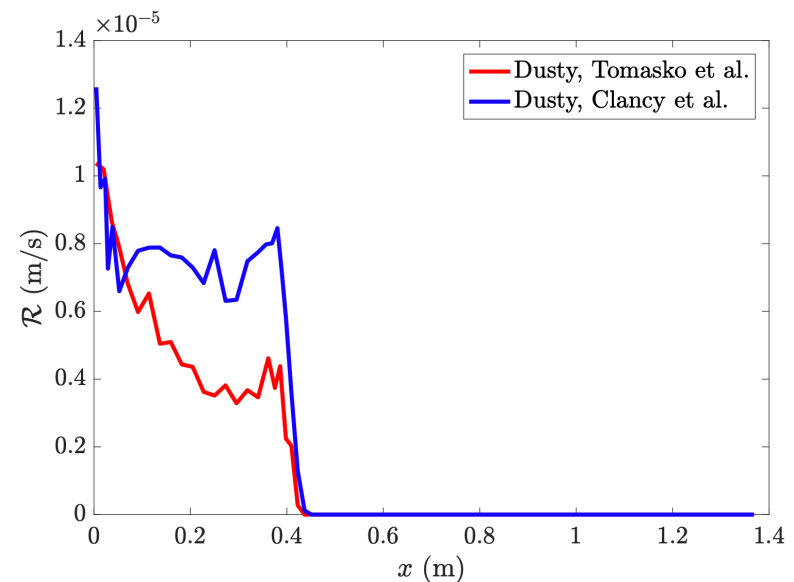
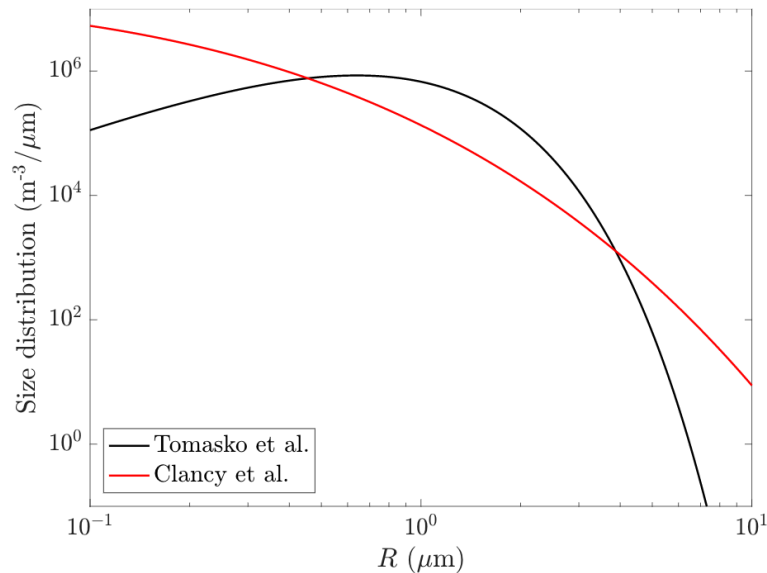
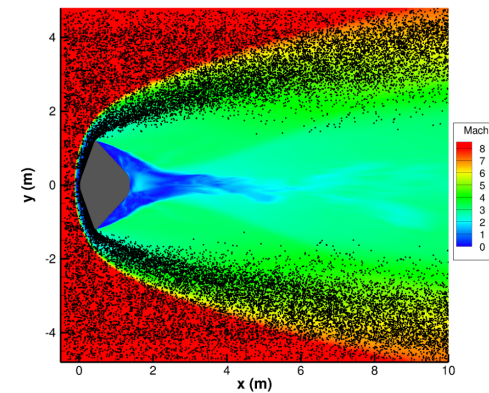
 $D = 0.5 \mu\text{m}$  $D = 2.0 \mu\text{m}$  $D = 5.0 \mu\text{m}$  $D = 8.0 \mu\text{m}$ 

Effect of particle size distribution

Modified gamma distribution:

- Tomasko et al., *J Geophys Research: Planets*, 1999
- Clancy et al., *J Geophys Research: Planets*, 1995

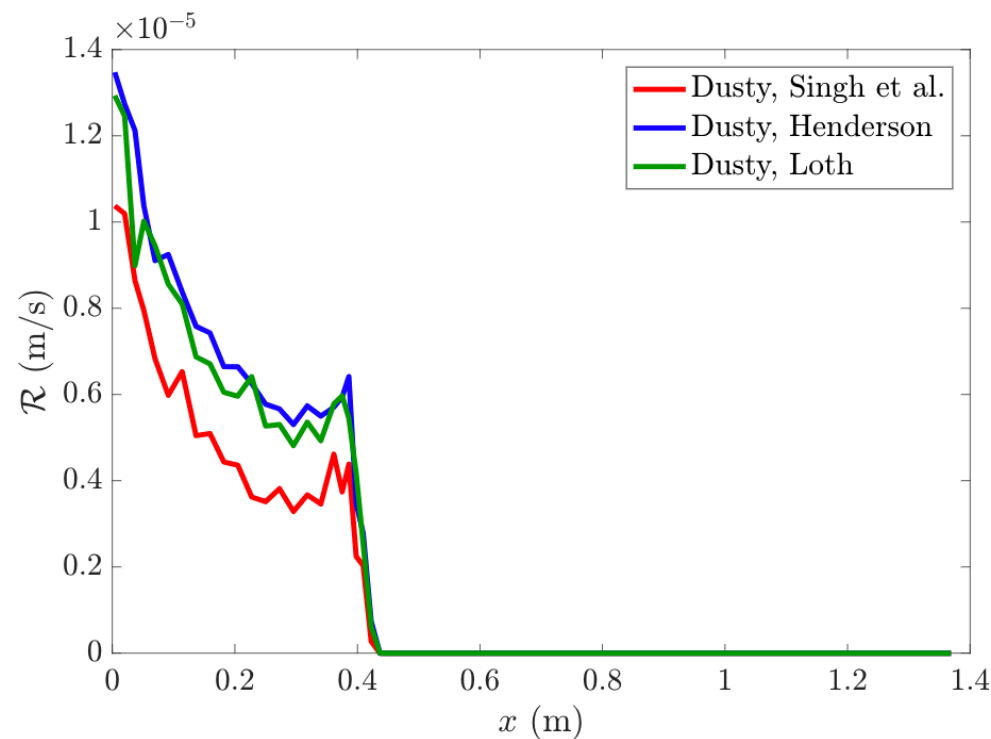
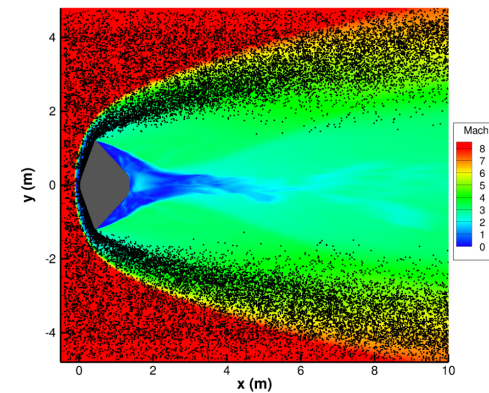
Mass loading ratio (particle mass flux to gas mass flux): $\beta = 0.01\%$



Effect of drag correlation

Drag correlations by:

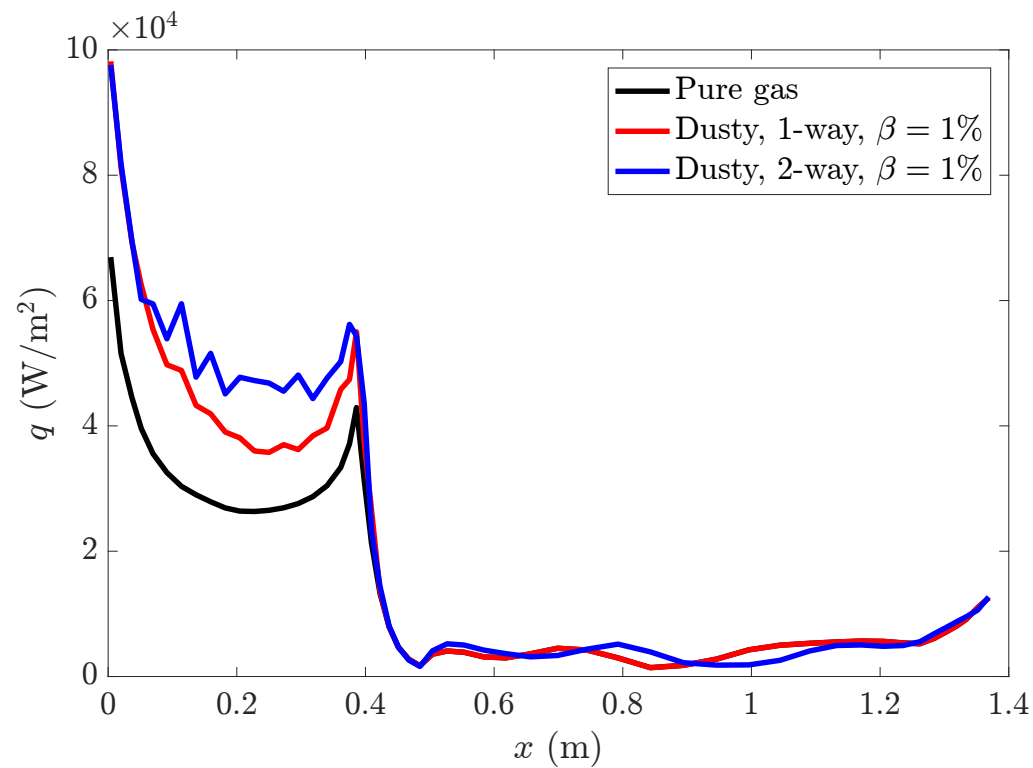
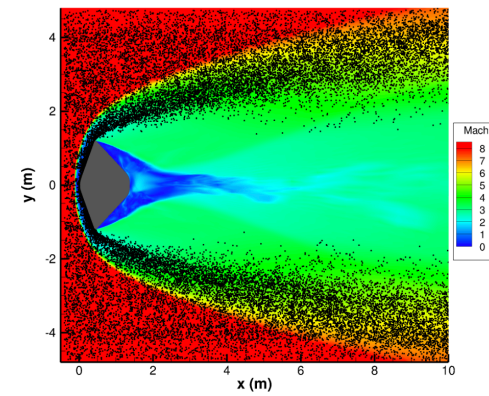
- Singh et al., <https://arxiv.org/pdf/2012.04813>
- Henderson, *AIAA J*, 1976
- Loth, *AIAA J*, 2008



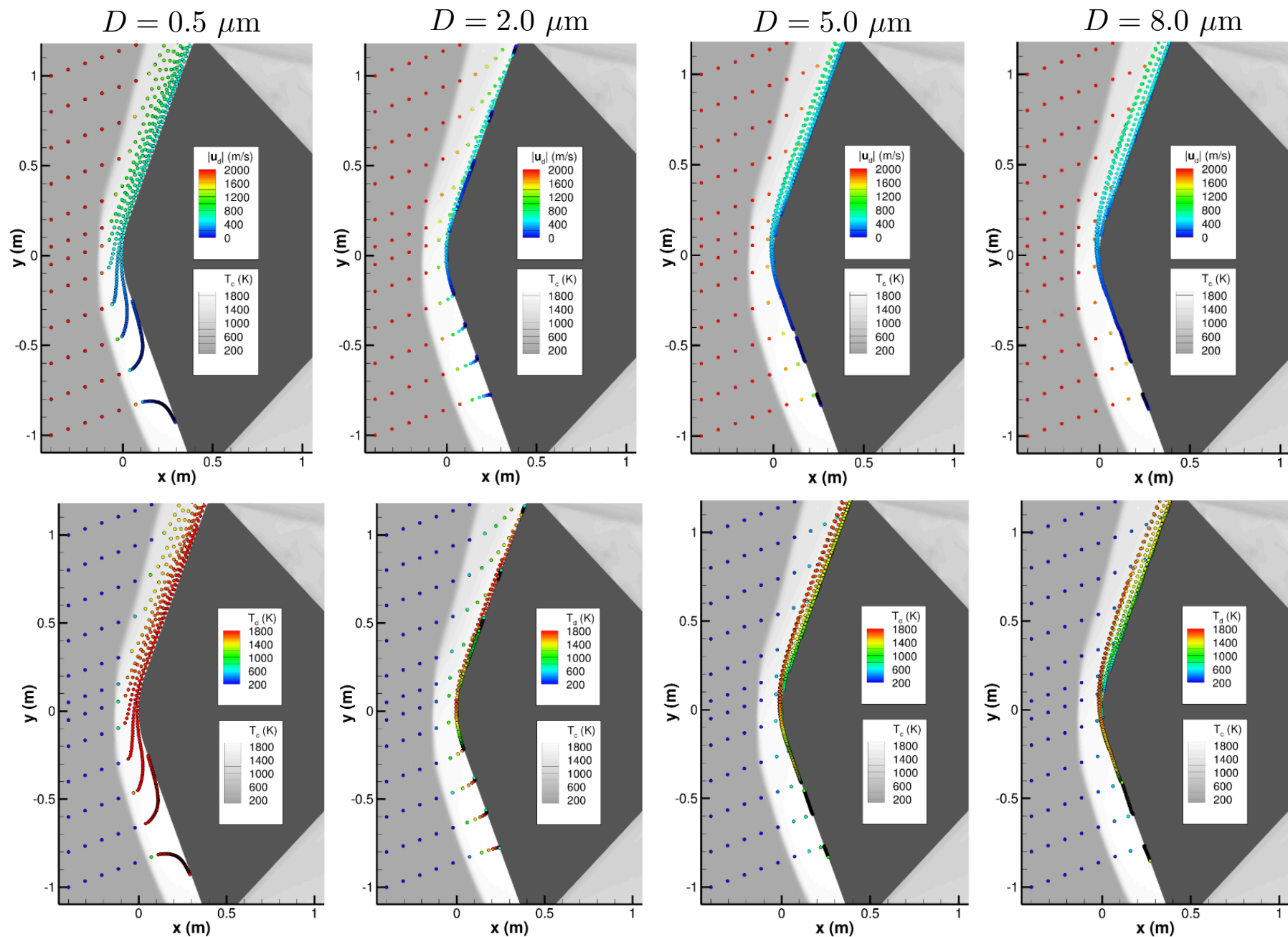
Effect of two-way coupling

Particles also affect carrier gas

$\beta = 1\%$ mass loading ratio



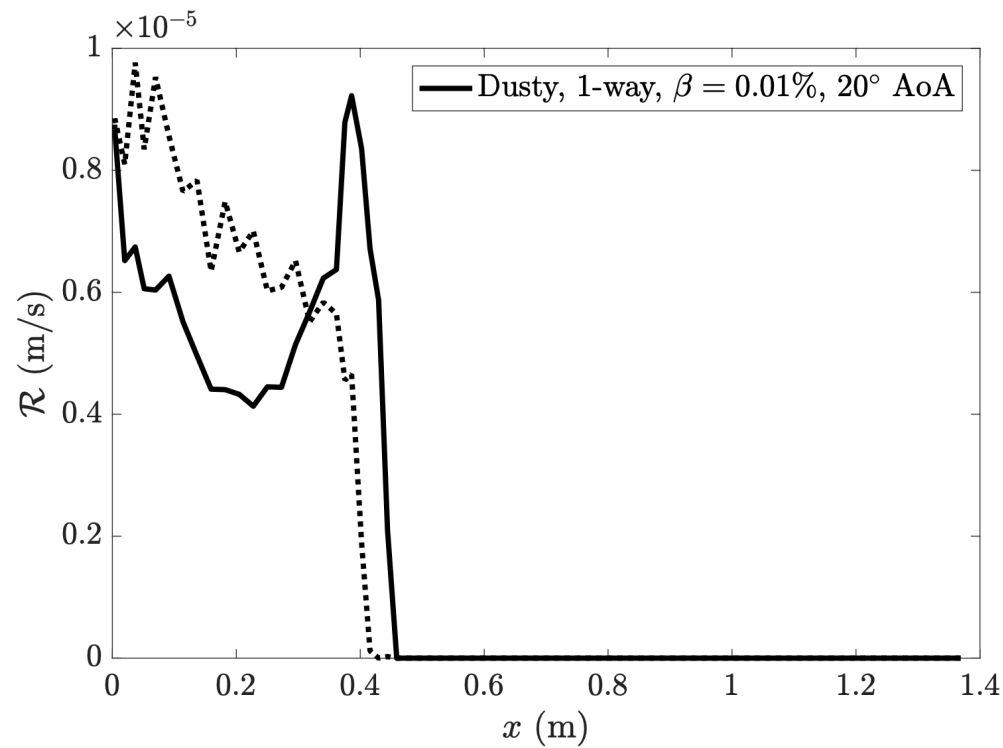
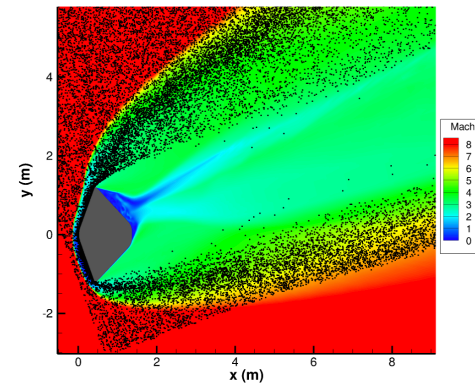
Sample particle trajectories (20 deg AoA)



Surface recession

Solid line: windward side

Dotted line: leeward side



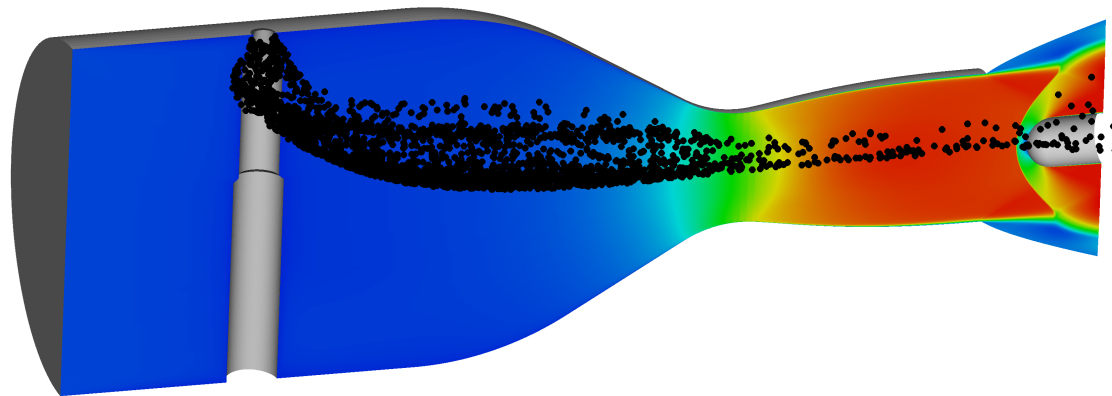
Main findings

- Dust impacts can cause surface erosion and heating augmentation
- Drag correlation by Singh et al. predicts less surface recession than popular correlations
- No direct interaction between particles and wake
- Large particles traverse the shock layer at higher speeds, lower temperatures than small particles
- Particle reverse-coupling at high dust loading can cause additional heating augmentation by transferring thermal energy to the boundary layer
- At nonzero angle of attack, leeward-side surface recession can be noticeably higher than windward-side recession

Summary & Outlook

Summary

- Developed a simple and robust shock capturing method for DG schemes
- Developed an Euler-Lagrange methodology in a DG framework suited for arbitrary curved, high-aspect-ratio elements
- Established significant benefits of the proposed DG methodology in the context of viscous hypersonic flows and particle-laden flows
- Applied the proposed methodology to hypersonic dusty flows over blunt bodies, with special application to Mars atmospheric entry



Outlook

- Further experimental validation: collaboration with DLR and NASA Ames
- Extend temperature range: thermochemical nonequilibrium

Thank you!

QUESTIONS

